AN ENERGY BARRIER IS NECESSARY FOR THE THERMAL STABILITY OF STABILIZER QUANTUM MEMORIES

by

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arXiv:1412.2858 arXiv: 1601.01324 THERMALIZATION TIME BOUND

• For a variety of models (commuting Pauli stabilizers and Abelian quantum doubles):

 $t_{Qmem} < t_{mix} \le \mathcal{O}(N^2 e^{2\beta\bar{\epsilon}})$

K. Temme. *arXiv:1412.2858* A. Kómár, O. Landon-Cardinal, K. Temme. *arXiv: 1601.01324*

OUTLINE

- Motivation
- Framework
- Thermalization time bounds
- Evaluation of the energy barrier
- Sketch of the proof



QUANTUM MEMORIES

• Quantum memory: physical system that encodes quantum information

- Degenerate ground space
- Robust to perturbations
- Long memory time (even in a thermal environment)
- Active error correction:





* Brown et al., *PRL (2014)*

THERMALIZATION IN 2D TORIC CODE

R. Alicki, M. Fannes, M. Horodecki J. Phys. A: Math. Theor. 42 (2009) 065303

- Thermalization represented by Liouvillian; with Gibbs state as a fixed point
- o Liouvillian: detailed balanced, with gap λ
- Mixing time:
 - State is close to Gibbs state
 - Bounded by N/λ

• Spectral gap bound: $\lambda \ge \frac{1}{3}e^{-8\beta J}$

• Implies mixing time bound: $t_{mix} \leq \mathcal{O}(Ne^{8\beta J})$

QUESTIONS

• Can we have a more rigorous connection between the energy barrier and thermalization?

• Is entropic protection possible?

FRAMEWORK

STABILIZER CODES

• Hamiltonian: $H = -J \sum_{k} g_{k}$

- With commuting Paulis: $\mathcal{G} = \{g_1, \dots, g_M\} \quad [g_i, g_j] = 0$
- Stabilizer group: $S = \langle G \rangle$ with codespace C(S) such that:
- $s \ket{\psi} = \ket{\psi} \; \; ext{for} \; \; egin{smallmatrix} s \in \mathcal{S} \ \ket{\psi} \in \mathcal{C}(\mathcal{S}) \end{cases}$
- Logical operators: $Comm(\mathcal{S}) \backslash \mathcal{S}$



 $\mathcal{H} = C^{2^N}$

ABELIAN QUANTUM DOUBLES

Kitaev "Fault-tolerant quantum computation by anyons." Annals of Physics (2003)

Generalization of Toric Code to Z_d qudits
Generalized Paulis:

$$X^{m} |k\rangle = |k \oplus_{d} m\rangle$$
$$Z^{l} |k\rangle = e^{\frac{2\pi i}{d} lk} |k\rangle$$



Source:

• Hamiltonian:



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NOISE MODEL



• Evolution: $\rho_S(t + \Delta t) = \operatorname{tr}[e^{-iH\Delta t}(\rho_S(t) \otimes \rho_R)e^{iH\Delta t}]$

• Markovian & Weak-Coupling limit:

$$\partial_t \rho_S = \mathcal{L}(\rho_S)$$

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DAVIES GENERATOR

E.B. Davies "Generators of dynamical semigroups." Journal of Functional Analysis (1979)

$$\mathcal{L}(\rho) = \sum_{\alpha,\omega} h^{\alpha}(\omega) \left(S_{\alpha}(\omega) \rho S_{\alpha}^{\dagger}(\omega) - \frac{1}{2} \{ S_{\alpha}^{\dagger}(\omega) S_{\alpha}(\omega), \rho \}_{+} \right)$$

• Jump operators:
$$e^{iH_S t} S_{\alpha} e^{-iH_S t} = \sum_{\omega} S_{\alpha}(\omega) e^{i\omega t}$$

• KMS condition:

$$h^{\alpha}(-\omega) = e^{-\beta\omega}h^{\alpha}(\omega)$$

• Fixed point: Gibbs state:

$$\rho_G \propto e^{-\beta H_S}$$

THERMALIZATION BOUNDS AND ENERGY BARRIER

MIXING TIME BOUND

• Mixing time: state is close to Gibbs state (trace-norm)

Theorem For any commuting N-qubit Pauli or N-qudit Abelian quantum double Hamiltonian, the mixing time is bounded by

$$t_{mix} < \mathcal{O}\left(N\frac{4\mu^*}{h^*}e^{2\beta\bar{\epsilon}}\right) \tag{1}$$

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- Largest (generalized) Pauli path: $\mu^* \leq \mathcal{O}(N)$
- Smallest transition rate: $h^* > c e^{-\beta \Delta}$
- Generalized energy barrier: ϵ

Pauli stabilizers: *Temme. arXiv:1412.2858* Abelian Quantum Doubles: *AK, Landon-Cardinal, Temme. arXiv: 1601.01324*

• Choose any error η_f (doesn't have to be a logical operator)



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- Throw out violated plaquettes (vertices)



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- Construct η_f by applying local errors
 Count number of violated plaquettes at each step



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- Take maximum of these for given ordering;
- But: choose an optimal ordering: $\bar{\epsilon}(\eta_f)$

$$\Rightarrow \bar{\epsilon} = \max_{\eta_f} \bar{\epsilon}(\eta_f)$$

EVALUATING THE ENERGY BARRIER FOR QUANTUM DOUBLES

- \circ (*d*-1) electric (magnetic) excitations + vacuum
- Break errors down to simple structures
- Generalized Pauli η_f = string-net of excitations





(d = 5)

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• String-net







EVALUATING THE ENERGY BARRIER FOR QUANTUM DOUBLES

• Helmholtz-Hodge decomposition:



• In general: η_f configuration = loops + trees

• Energy barrier:

- Of loops: 2
- Of trees (simple union of strings): 2

Sketch of Proof

SPECTRAL GAP AND THERMALIZATION TIME

• Convergence to unique fixed point σ :

$$||e^{t\mathcal{L}}(\rho_0) - \sigma||_{tr} \le \sqrt{||\sigma^{-1}||}e^{-\lambda t}$$

• For thermal σ :

$$||\sigma^{-1}|| \sim e^{c\beta N} \qquad \Longrightarrow \qquad t_{mix} \sim \mathcal{O}(\beta N \lambda^{-1})$$

POINCARÉ INEQUALITY

$$\lambda = \min_{f \in \mathcal{M}_{d^N}} \frac{\mathcal{E}(f, f)}{\operatorname{Var}_{\sigma}(f, f)}$$

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POINCARÉ INEQUALITY

$$\lambda = \min_{f \in \mathcal{M}_{d^N}} \frac{-\operatorname{tr}[\sigma f^{\dagger} \mathcal{L}^*(f)]}{\operatorname{tr}[\sigma f^{\dagger} f] - \operatorname{tr}[\sigma f]^2}$$

• Any $\mu > 0$ such that $\mu \le \frac{\mathcal{E}(f, f)}{\operatorname{Var}_{\sigma}(f, f)}$ for all f is a lower bound to $\lambda \implies \operatorname{maximize} \mu$

• Equivalently: minimize τ subject to $\tau \hat{\mathcal{E}} - \hat{\mathcal{V}} \ge 0$ $\mathcal{E}(f, f) = (f|\hat{\mathcal{E}}|f)$ $\operatorname{Var}_{\sigma}(f, f) = (f|\hat{\mathcal{V}}|f)$

(QUANTUM) CANONICAL PATHS

• With defining a quantum canonical path we can factorize: $\hat{\mathcal{E}} = AA^{\dagger}$ and $\hat{\mathcal{V}} = BB^{\dagger}$ and define AW = B

• Then the spectral gap bound: $\tau = \min_{W} ||W||^2$

• For classical Markov processes:

- M. Jerrum, A. Sinclair. "Approximating the permanent." SIAM journal on computing 18.6 (1989): 1149-1178.
- Geometric picture: graph of states

• Generalized to Quantum setting:

• K. Temme. "Thermalization time bounds for Pauli stabilizer Hamiltonians." *arXiv preprint arXiv:1412.2858* (2014).

CONCLUSION & OPEN QUESTIONS

- Commuting Pauli stabilizer Hamiltonians and Abelian quantum doubles follow the Arrhenius law
 - Even with defect lines of Brown et al., PRL (2014) $t_{mix} \leq \mathcal{O}(\beta N^2 e^{2\beta \bar{\epsilon}})$
- No self-correction without a scaling energy barrier
- Not excluded of entropic protection:
 - Different type of defects (*Bombin*, *PRL* (2010))
 - Different noise model
 - Non-Abelian quantum doubles

K. Temme. *arXiv:1412.2858* A. Kómár, O. Landon-Cardinal, K. Temme. *arXiv: 1601.01324*



• Path on the Pauli group: $\eta_0, \eta_1, \ldots \eta_f \in \mathcal{P}_N$

GENERALIZED ENERGY BARRIER

COMMUTING PAULI STABILIZERS

• Reduced set \mathcal{G}_{η} :

- Removing stabilizer generators violated in the final error configuration, we only consider the other stabilizer generators
- Energy barrier of a Pauli:

$$\bar{\epsilon}(\eta_f) = \min_{\eta_t} \max_t 2J^* \# \{ g_k \in \mathcal{G}_\eta \mid [g_k, \eta_t] \neq 0 \}$$

• Generalized energy barrier: $\bar{\epsilon} = \max_{\eta_f} \bar{\epsilon}(\eta_f)$ √ •**Z**-----O-----O-----O-----O---

Defined for

any operator,

not only logical ops.



 $t_{mix} < \beta 8N^2 e^{\beta 6J}$

For 2D Toric Code:

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GENERALIZED ENERGY BARRIER – Abelian Quantum Doubles

Defined for **any** operator, not only logical ops.

- Path on the *generalized* Pauli group: $\eta_0, \eta_1, \ldots, \eta_f \in \mathcal{P}_N$
- Set of projectors: $\mathcal{A} = \{P_v^l, Q_p^l \mid l = 0, 1, \dots d 1, all v, p\}$
- Reduced set \mathcal{A}_{η} :
 - Removing operators violated in the final error configuration, we only consider the other elements of \mathcal{A}
- Energy barrier of a Pauli:

 $\bar{\epsilon}(\eta_f) = \min_{\eta_t} \max_t 2J^* \#\{a_k^l \in \mathcal{A}_\eta \mid [a_k^l, \eta_t] \neq 0\}$

• Generalized energy barrier:

$$\bar{\epsilon} = \max_{\eta_f} \bar{\epsilon}(\eta_f)$$

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