Information propagation for interacting particle systems

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joint work with Sarah Harrison, Tobias Osborne, and Jens Eisert

Introduction

• How fast can **information propagate** in physical systems?

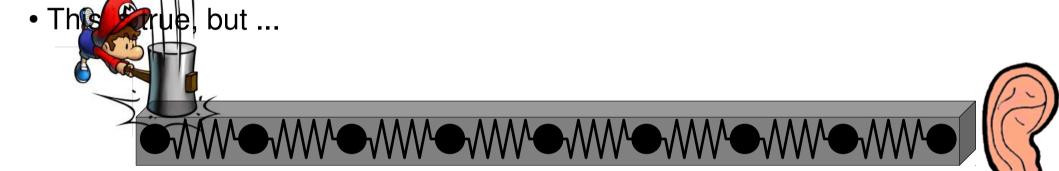
→ Obvious answer: Relativity ⇒ No faster than the speed of light!

• This is true, but ...



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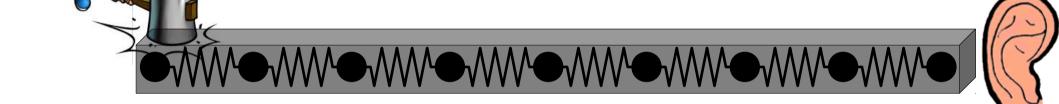


... e.g. in classical mechanical systems, information propagates at a **speed of sound**, without the need for relativistic arguments!

- This speed can be understood from the microscopic model, using
 - that it is local
 - that the interactions have bounded strength

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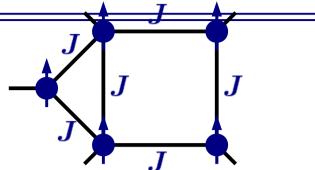
- that the interactions have bounded strength

 \Rightarrow Finite propagation speed can be understood non-relativistically!

Quantum mechanical systems

- What about quantum mechanical systems?
- Quantum spin systems:

$$- \downarrow J \downarrow J \downarrow J \downarrow J \downarrow$$



 $H = \sum_{\langle j,k \rangle} h_{jk}$; $\|h_{jk}\|_{op} \leq J$: local Hamiltonian of bounded strength

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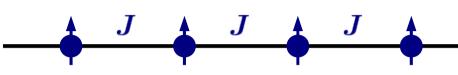
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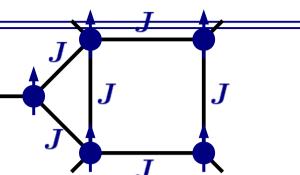
• Lieb-Robinson bounds: [Lieb & Robinson '72, Hastings '04, Nachtergaele & Sims '06] $\|[A(t), B]\| \le c \|A\| \|B\| \exp[-(L - vt)/\xi]$ Lieb-Robinson velocity $v = c_G J$

depends on graph

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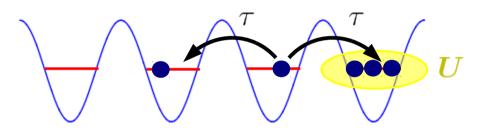
• Lieb-Robinson bounds: $\|[A(t), B]\| \le c \|A\| \|B\| \exp[-(L - vt)/\xi]$ Lieb-Robinson velocity $v = c_G J$ depends on graph • Relevance: • question of fundamental interest • propagation speed of perturbations • facilitates simulation of dynamics • imaginary time \Rightarrow exponential decay of correlations • What about systems of interacting particles, such as bosons?

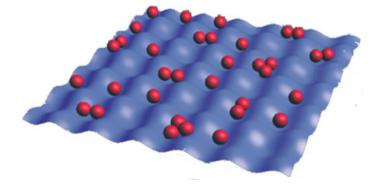
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• canonical example: **Bose-Hubbard model:**

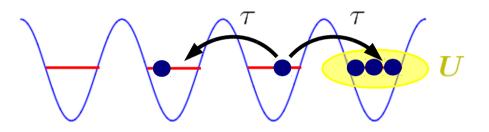


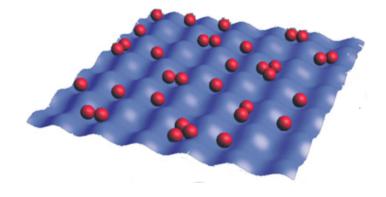


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canonical example: Bose-Hubbard model:





$$H_{\rm BH} = -\tau \sum_{\langle j,k \rangle} (\hat{a}_j^{\dagger} \hat{a}_k + \hat{a}_k^{\dagger} \hat{a}_j) + U \sum_j \hat{n}_j (\hat{n}_j - 1)$$

 \hat{a}_{j} : annihilate a particle at site j \hat{a}_{j}^{\dagger} : create a particle at site j $\hat{a}_{j}|n\rangle = \sqrt{n}|n-1\rangle \quad \leftrightarrow \quad \hat{n}_{j} = \hat{a}_{j}^{\dagger}\hat{a}_{j}$: counts particles at site

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- Lieb-Robinson bound does not apply:

$$a_j^{\dagger} a_k |n_j - 1, n_k\rangle = \sqrt{n_j n_k} |n_j, n_k - 1\rangle$$

 \Rightarrow hopping term $a_j^{\dagger}a_k$ unbounded (or only by $||a_j^{\dagger}a_k|| \leq N_{\text{tot}}$)

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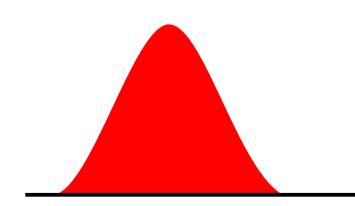
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 - \Rightarrow need **constraints on Hamiltonian** (e.g. particle number conserving)
- hopping rate (and thus v) will depend on the filling of the lattice:
 ⇒ need constraints on initial state
- Note: bounds exist for quadr. Hamiltonians and certain perturbations thereof [Nachtergaele, Raz, Schlein, Sims 2009]

Idea: Restrict to relevant models

- Aim: propagation speed for **Bose-Hubbard type models**
- How can we obtain a meaningful propagation speed?
 - restrict to certain initial states of interest (which allow for finite speed of propagation)
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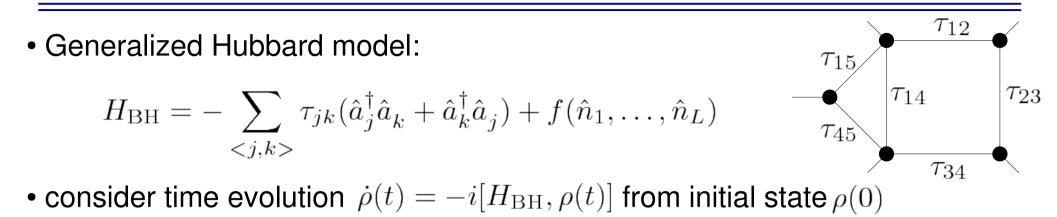
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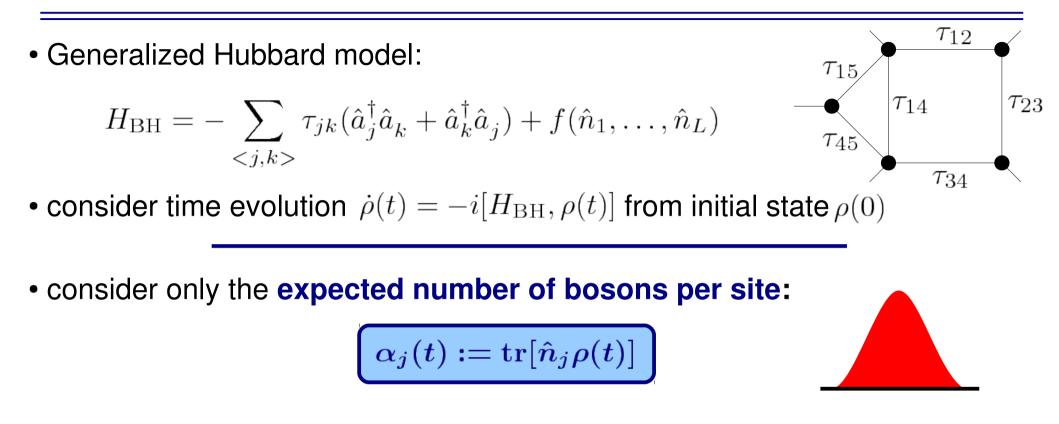
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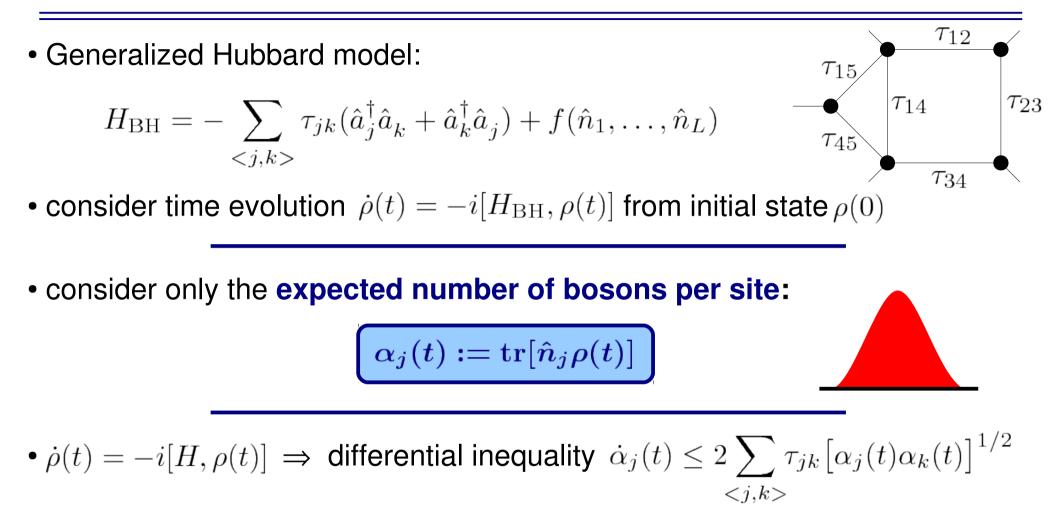
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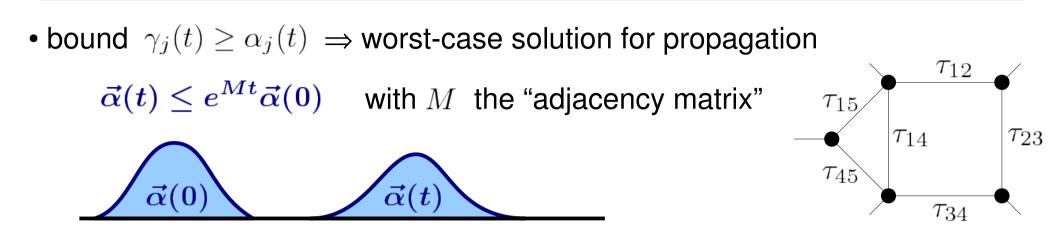
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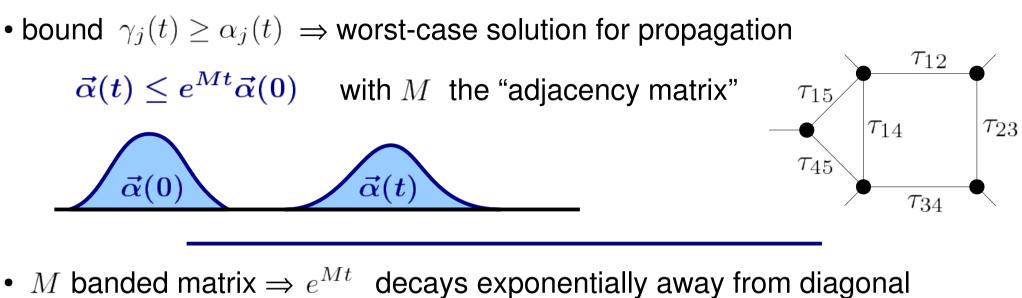
 \Rightarrow worst-case upper bound $\gamma_j(t) \ge \alpha_j(t)$ evolves according to:

$$\dot{\gamma}_j(t) = 2 \sum_{\langle j,k \rangle} \tau_{jk} (\gamma_j(t) + \gamma_k(t))$$
 (linearized)

Obtaining a speed limit



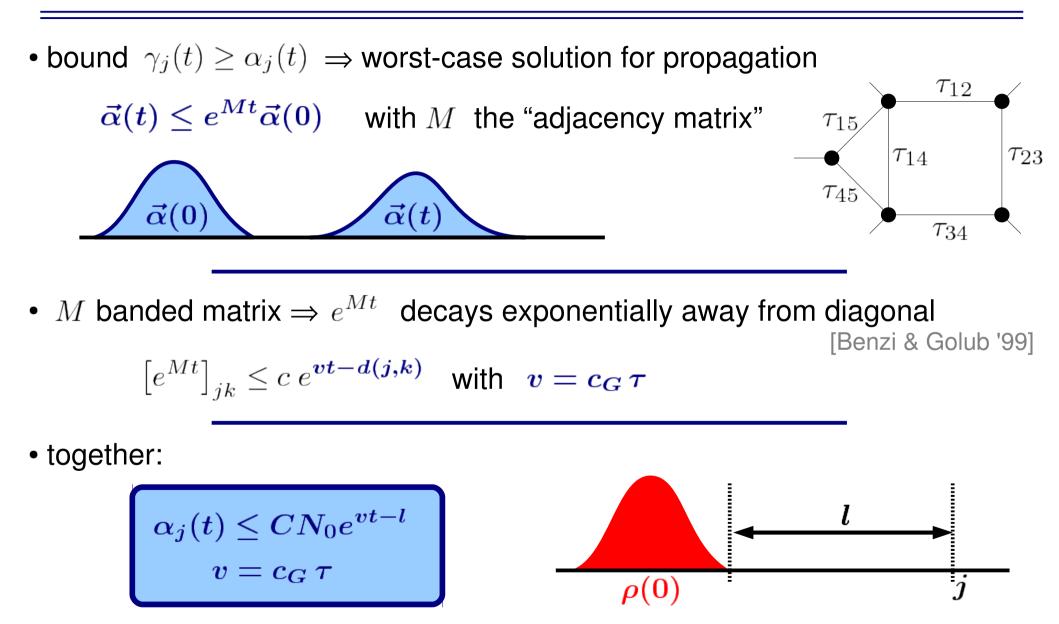
Obtaining a speed limit



[Benzi & Golub '99]

$$\left[e^{Mt}
ight]_{jk} \leq c \; e^{oldsymbol{vt} - oldsymbol{d}(oldsymbol{j},oldsymbol{k})} \quad ext{with} \; \; oldsymbol{v} = oldsymbol{c}_{oldsymbol{G}} \, oldsymbol{ au}$$

Obtaining a speed limit



 \Rightarrow speed independent of particle number!

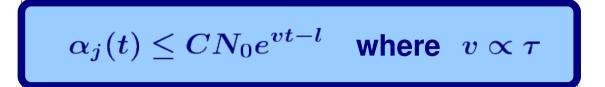
Speed limit for interacting particles

$$lpha_j(t) \leq C N_0 e^{vt-l}$$
 where $v \propto au$

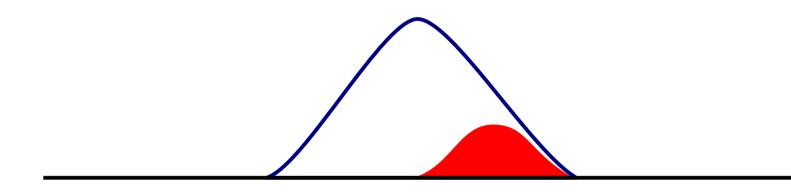
• Proof idea: Study evolution of worst case bound on $\alpha_j(t)$:



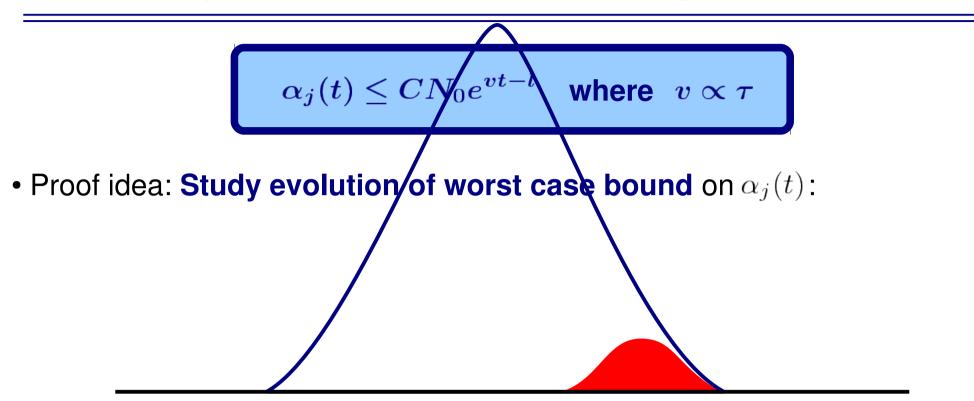
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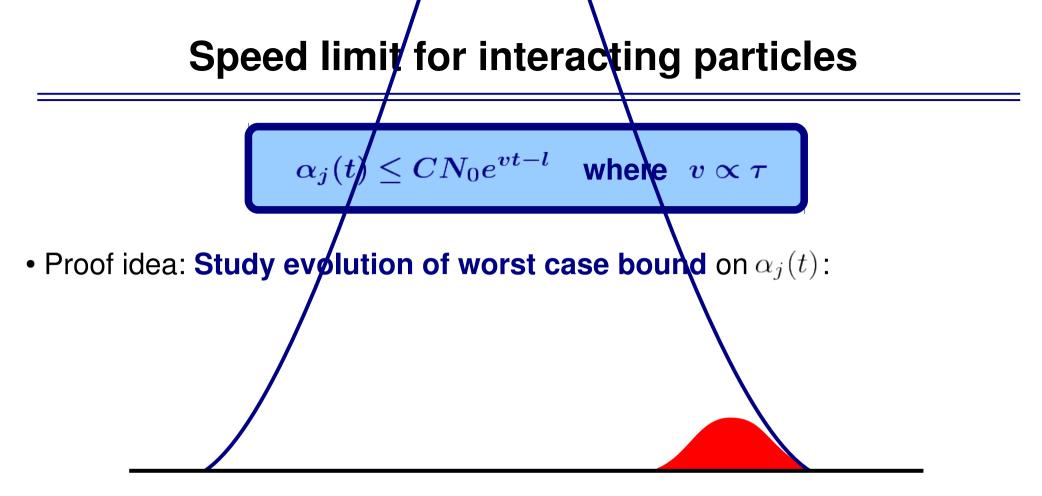


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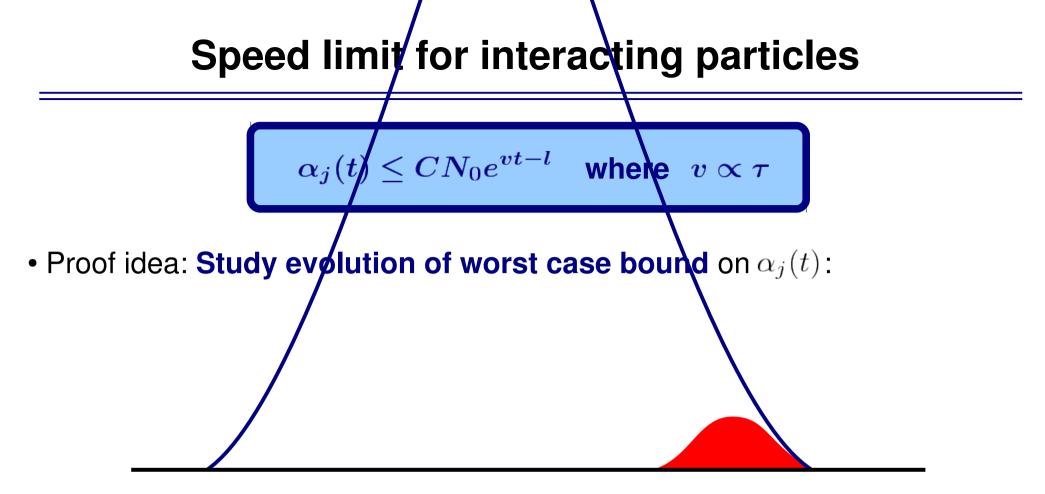


Speed limit for interacting particles





argument works for any Hubbard-type model on any graph



- argument works for any Hubbard-type model on any graph
- extension possible to
 - higher moments of particle number
 - arbitrary local operators
 - operators acting on larger blocks (up to log-size)

... by iteratively bounding those quantities by $lpha_j(t)$.

 can be extended to several species of particles, fermions, Bose-Fermi mixtures, and even anyons:

$$H = -\sum_{\langle j,k \rangle,s} \tau_{jk} (\hat{a}_{j,s}^{\dagger} \hat{a}_{k,s} + \hat{a}_{k,s}^{\dagger} \hat{a}_{j,s}) + f(\{n_{j,s}\}_{j,s})$$

 \rightarrow can be understood as **hopping on independent graphs**

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• works for certain dissipative theories, e.g. for particle losses:

$$\dot{\rho}(t) = -i[H_{BH}, \rho(t)] - \mathcal{L}[\rho(t)] \qquad \text{describes} \\ \text{loss of particles}$$

$$\Rightarrow \dot{\boldsymbol{\alpha}}_{\boldsymbol{j}}(\boldsymbol{t}) = \dot{\alpha}_{\boldsymbol{j}}^{\operatorname{Ham}}(t) - \underbrace{\dot{\alpha}_{\boldsymbol{j}}^{\operatorname{diss}}(t)}_{\geq 0} \leq \dot{\boldsymbol{\alpha}}_{\boldsymbol{j}}^{\operatorname{Ham}}(\boldsymbol{t})$$

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• idea extendible to **continuum theories**: either continuum limit, or continuous differential inequalities for $\alpha(x, t)$

Summary

- We have studied the propagation of interacting bosons
- We have found a finite propagation speed for any excitation into the initially unoccupied region
- Propagation speed only depends on coupling strength
- Extends to Bose-Fermi mixtures, dissipative models, continuum theories

arXiv:1010.4576