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# Anyons, Twists & Topological Codes

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# outline

- motivation/introduction
- twists
  - anyon models, symmetries & twists
  - toric code
- topological subsystem color codes
  - twist geometry
  - Clifford operations by code deformation
- conclusions

# motivation

- traditional fault-tolerance is not enough to reallistically fight decoherence
- topological alternatives (Kitaev):
  - topological quantum computing (TQC)
    - self-protected by energy gap
    - immune to small local distortion
  - topological codes (TC)
    - geometrically local, active error correction
    - error threshold for large size
- both are anyon-based: exotic statistics in 2D

### motivation

- two problems addressed here:
  - TQC: the anyons that are easier to get have no computational power
  - TC: there exist extremely local TCs (2-local measurements in 2D), but no way to compute
- a solution / new tool:
  - twists → use anyon symmetries to increase computational power

• why 2D?



- statistics beyond bosons and fermions
- topological interaction
- appear in systems with topological order (TO) (Wen '89):
  - gapped, ground state degeneracy depends on topology

- abelian charge: given charge of constituents, total charge is known
- topological charge can be **non-abelian**



• TQC (Kitaev '03, Freedman et al '03)

encode in fusion channels



t



measure = fuse





- abelian anyons have no computational power
- twists offer a way to recover non-abelian behavior!

- quantum error correcting codes protect quantum information using redundancy
- typically this involves encoding in a subspace

Hilbert space



 the code subspace can be defined in terms of commuting observables: check operators (CO)

$$C_i|\psi\rangle = c_i|\psi\rangle$$

 errors typically change CO values → allows to keep track of errors

CO measurement  $\rightarrow$  error syndrome  $\rightarrow$ 

 $\rightarrow$  compute most probable error

- topological codes (Kitaev '97)
  - geometrically **local** check operators = easy to measure
  - global undetectable errors = hard to happen



- # encoded qubits depends on topology (homology)
- flexible: many lattices allowed, transversal gates possible
- **boundaries**: planar geometries
- topological quantum memory (Dennis et al '02):
  - measure COs repeatedly
  - under a noise **threshold**, storage time exp in size
  - ideal error correction amounts to compute free energy
- code deformation:
  - change topology over time: initialize, compute, measure

- subsystem codes (Kribs et al '05) can also improve locality
- only a subsystem of the code subspace is used
- check operators need not be measured directly → measurements potentially more local (Poulin '05)



- **topological subsystem color codes (TSCC)** (Bombin '09)
  - "doubly local": topology + subsystem
  - error syndrome recovery needs 2-local measurements!



- TC and TO are closely related for subspace codes
  - Topological codes
  - Check operators
    - Code subspace
    - Error syndrome

- VS <u>Topological order</u>
- $\leftrightarrow$  Hamiltonian terms
- ↔ Ground subspace
- ↔ Excitation configuration
- TSCCs also have an anyonic picture for error syndromes

- TSCCs do not allow boundaries
  - no natural planar codes
  - code deformation becomes unpractical
- with twists
  - we can build planar TSCCs
  - whole Clifford group by code deformation!



ingredients of an anyon model:









 $|\psi\rangle \longrightarrow ?$ 

topological charges

fusion rules

braiding rules

#### • ex.: Ising anyons

- topological charges Q {1,  $\sigma$ ,  $\psi$ }
- fusion rules

 $\sigma \times \sigma = \mathbf{1} + \psi, \qquad \sigma \times \psi = \sigma, \qquad \psi \times \psi = \mathbf{1}.$ 

- the total charge of two distant  $\sigma$ -s is 1 or  $\psi$ :
  - if far appart, global qubit
- fusion space:  $2n \sigma$ -s  $\rightarrow n$  qubits

- braiding rules:
  - we can describe braiding up to a phase with a Majorana operator per  $\sigma$ 
    - Majorana operators are self-adjoint  $c_i$  with

$$c_j c_k + c_k c_j = 2\delta_{jk}$$

• total charge of *j*-th and *j*+1-th  $\sigma$ -s: (Q)



$$-ic_jc_{j+1}$$

- braiding:  $c_j \to c_{j+1}$   $\overbrace{c_{j+1}}^{i \to i} \to -c_j$
- not universal, but we can use distillation (Bravyi '06)

 anyon symmetry: charge permutation producing an equivalent anyon model

$$q \longrightarrow \pi(q)$$

 imagine 'cutting' the anyons' 2D world and gluing it again up to a symmetry

$$\longrightarrow \longrightarrow \longrightarrow$$

• across the cut, charges change:

$$\begin{array}{c} q \\ \longrightarrow \\ \bullet \end{array} \\ \bullet \end{array} \\ \bullet \\ \mathbf{x} \end{array}$$

- topologically, the cut location is unphysical.
- endpoints are meaningful: under monodromy they permute charges → twists



- ex.: quantum double of Z<sub>2</sub> (toric code)
- charges:  $\{\mathbf{1}, e, m, \epsilon\}$
- fusion:  $e \times m = \epsilon$   $e \times \epsilon = m$   $m \times \epsilon = e$  $e \times e = m \times m = \epsilon \times \epsilon = 1$
- braiding:  $e, m \rightarrow bosons$   $\epsilon \rightarrow fermion$

$$\bigvee_{q} q' = - \bigcap_{q} \left( q' \quad \mathbf{1} \neq q \neq q' \neq \mathbf{1} \right)$$

• nontrivial symmetry:  $e \leftrightarrow m$ 

• twists are sinks/sources for fermions:



vacuum to vacuum processes...

$$(\mathbf{x}) \in \mathbf{x}$$

• ...lead to **topological degeneracy**:



- toric code (Kitaev '97, Wen '03):
  - qubits form a square lattice
  - 4-local check operators at plaquettes



$$C_k := X_k Z_{k+\mathbf{i}} X_{k+\mathbf{i}+\mathbf{j}} Z_{k+\mathbf{j}}$$
$$C_k |\psi\rangle = |\psi\rangle$$

Hamiltonian version:

$$H := -\sum_{k} C_{k}$$

• excitations live at plaquettes

 string operators create/destroy excitations at their endpoints



• two types of strings/excitations: e (light) and m (dark)

• twists amount to **dislocations** 



• twists can be locally created in PAIRS only

• no twists (or even number)  $\rightarrow$  4 possible charges



• a twist (or an odd number)  $\rightarrow$  2 possible charges



• non-abelian fusion rules!

$$\sigma_{\pm} \times \sigma_{\pm} = 1 + \epsilon \qquad \sigma_{\pm} \times \sigma_{\mp} = e + m$$
  
$$\sigma_{\pm} \times \epsilon = \sigma_{\pm} \qquad \sigma_{\pm} \times e = \sigma_{\pm} \times m = \sigma_{\mp}$$

• we recover **Ising** rules:

 $\sigma_+ \times \sigma_+ = \mathbf{1} + \epsilon \qquad \sigma_+ \times \epsilon = \sigma \qquad \epsilon \times \epsilon = \mathbf{1}$ 

 all closed string ops can be expressed in terms of a set of open string ops → Majorana operators



$$c_j c_k + c_k c_j = 2\delta_{jk}$$

• braiding is also Ising-like!





- the original TSCCs come from 3-valent lattices with
   3-colorable faces (red, green, blue)
- string operators have a color
- commutation relations of string ops relates them to an anyon model with three nontrivial charges



• fusion rules as in toric code

$$r \times g = b$$
  $g \times b = r$   $b \times r = g$   
 $r \times r = g \times g = b \times b = 1$ 

- braiding of different charges as in toric code
- the difference: three fermionic charges
- any **permutation** of the colors is a symmetry!
- twists are labeled by the elements of  $S_3$



- faces with an odd number of links brake 3-colorability
- these are twists: two colors are exchanged
- a red twist exchanges green and blue, and so on



• to the i-th twist we attach a string  $\gamma_{i...}$ 



...and get self-adjoint string ops k<sub>i</sub>,
 "colored" Majorana ops

 $k_i^2 = 1 \text{ and, for } i < j,$   $k_i k_j = \begin{cases} k_j k_i & \text{if } c_i = \zeta_+(c_j), \\ -k_i k_j & \text{otherwise.} \end{cases}$ 

• braiding changes the color of twists



 transforming as follows the colored Majorana ops (c<sub>i</sub> is the color of the i-th twist)

$$k_j \to k_{j+1}, \ k_{j+1} \to \begin{cases} -k_j & \text{if } c_j = c_{j+1}, \\ ik_j k_{j+1} & \text{if } c_j = \zeta_-(c_{j+1}), \\ -k_j k_{j+1} & \text{otherwise.} \end{cases}$$

- for twists of the same color, we are back to Ising anyons
- encoding: I qubit = 4 twists of the same color
- we get all single qubit Clifford gates (Bravyi '06)

$$\langle k_j k_{j+1} k_{j+2} k_{j+3} \rangle = -1$$

$$\hat{X} \equiv -ik_j k_{j+1}$$

$$\hat{Z} \equiv -ik_{j+1} k_{j+2}$$

- to get the whole Clifford group, we only need to implement an entangling gate
- but for two groups of twists of different color:

$$\hat{X}_1 \to \hat{X}_1 \qquad \hat{Z}_1 \to \hat{X}_2 \hat{Z}_1$$
$$\hat{X}_2 \to \hat{X}_2 \qquad \hat{Z}_2 \to \hat{X}_1 \hat{Z}_2$$

and we can always flip the color of a group:



TSCC + twists + + code deformation = Clifford gates

# conclusions & questions

- anyon symmetries allow to introduce twists
- twists make anyon models and topological codes computationally more powerful (but how much?)
- toric codes:
  - twists mimic Ising anyons
- topological subsystem color codes:
  - Clifford operations by code deformation
- other/general anyon models?

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