

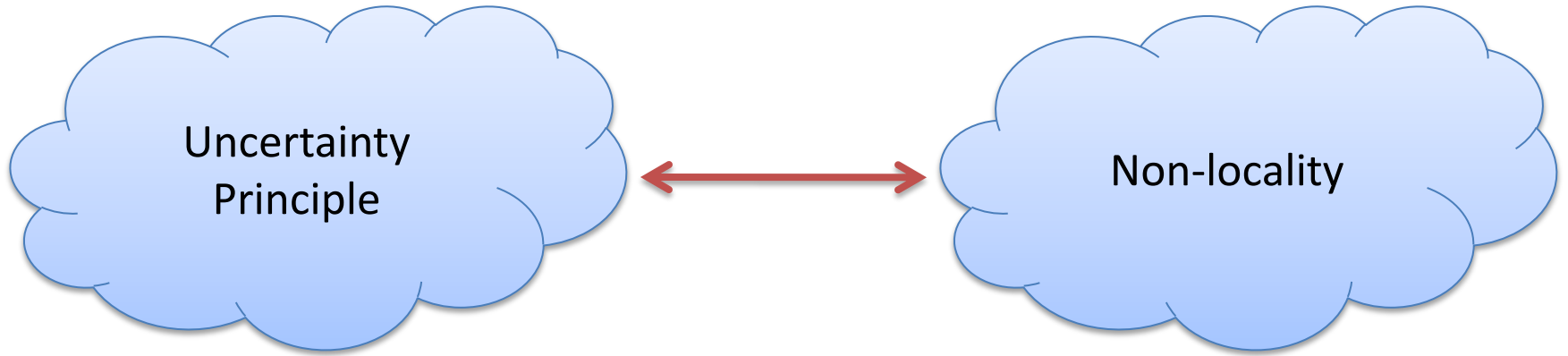
# The uncertainty principle determines the non-locality of Quantum Mechanics

Stephanie Wehner

Joint work with Jonathan Oppenheim  
Science, vol 330, no 7006, pp 1072-1074



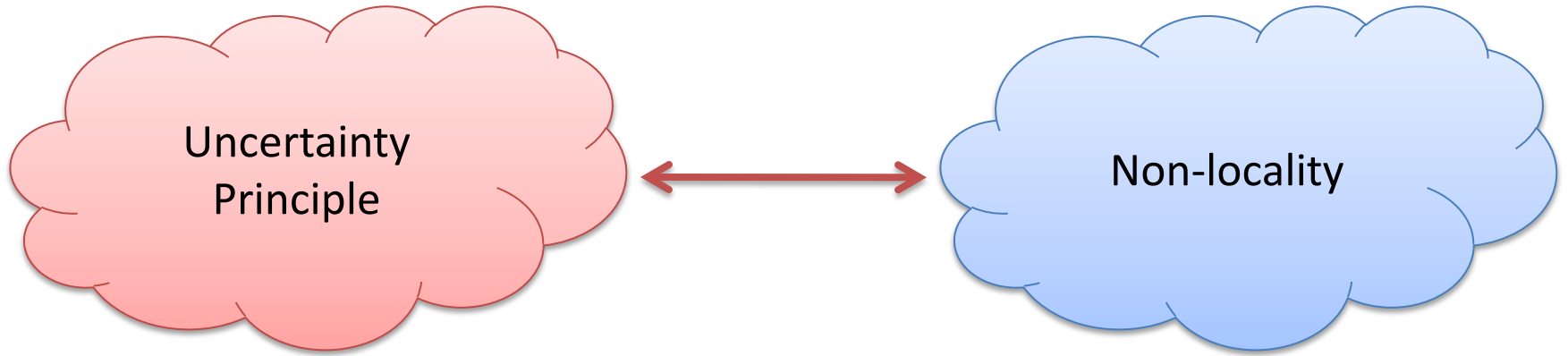
# Linking two fundamental concepts



Steering

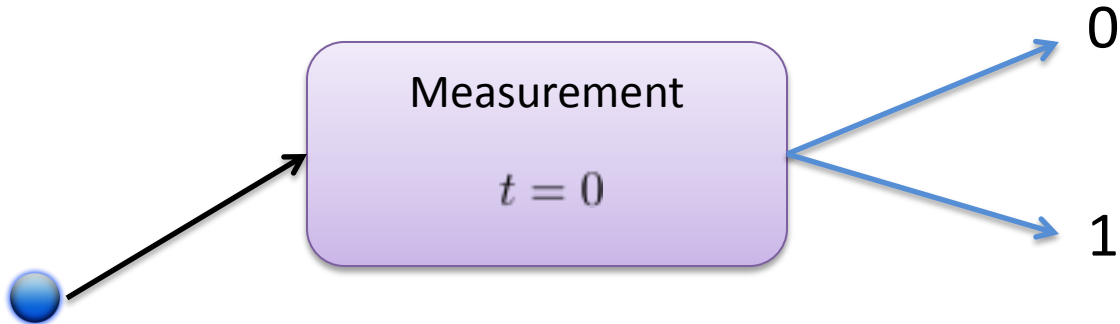


# Linking two fundamental concepts



Steering

# What is uncertainty?



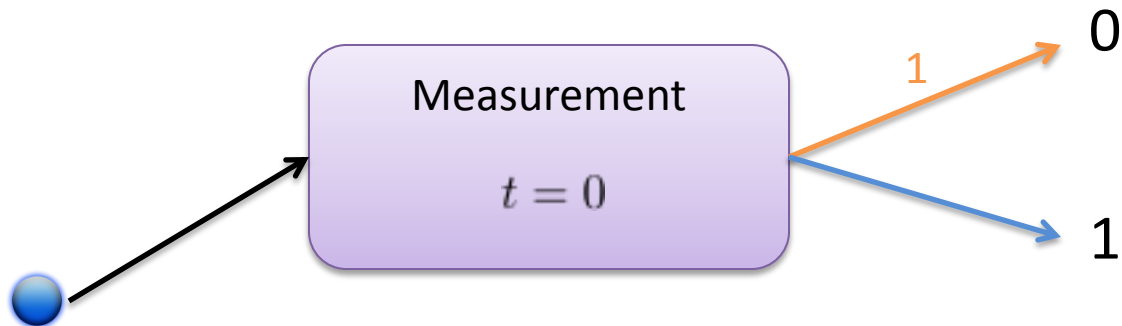
**$p(b | t)$**   
Probability of obtaining  
outcome bit **b**  
when performing  
measurement **t**

“Cannot know both the particle’s position and its momentum at the same time” observed by Werner Heisenberg in 1927.





# What is uncertainty?



**$p(b | t)$**

Probability of obtaining outcome bit **b**

when performing measurement **t**

Certainty:

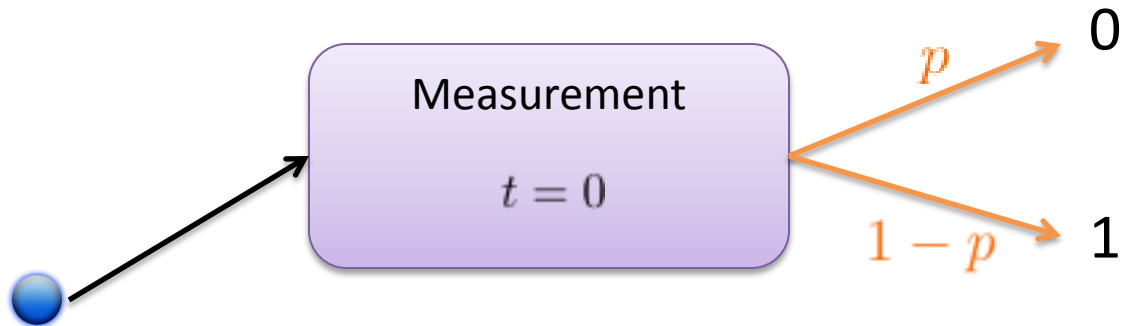
We know the property to be determined perfectly.

For some outcome  $b$  (here  $b = 0$  ) we have  $p(b|t) = 1$

“Cannot know both the particle’s position and its momentum at the same time” observed by Werner Heisenberg in 1927.



# What is uncertainty?



**$p(b | t)$**   
Probability of obtaining  
outcome bit **b**  
  
when performing  
measurement **t**

Uncertainty:

We do *not* know the property to be determined perfectly.

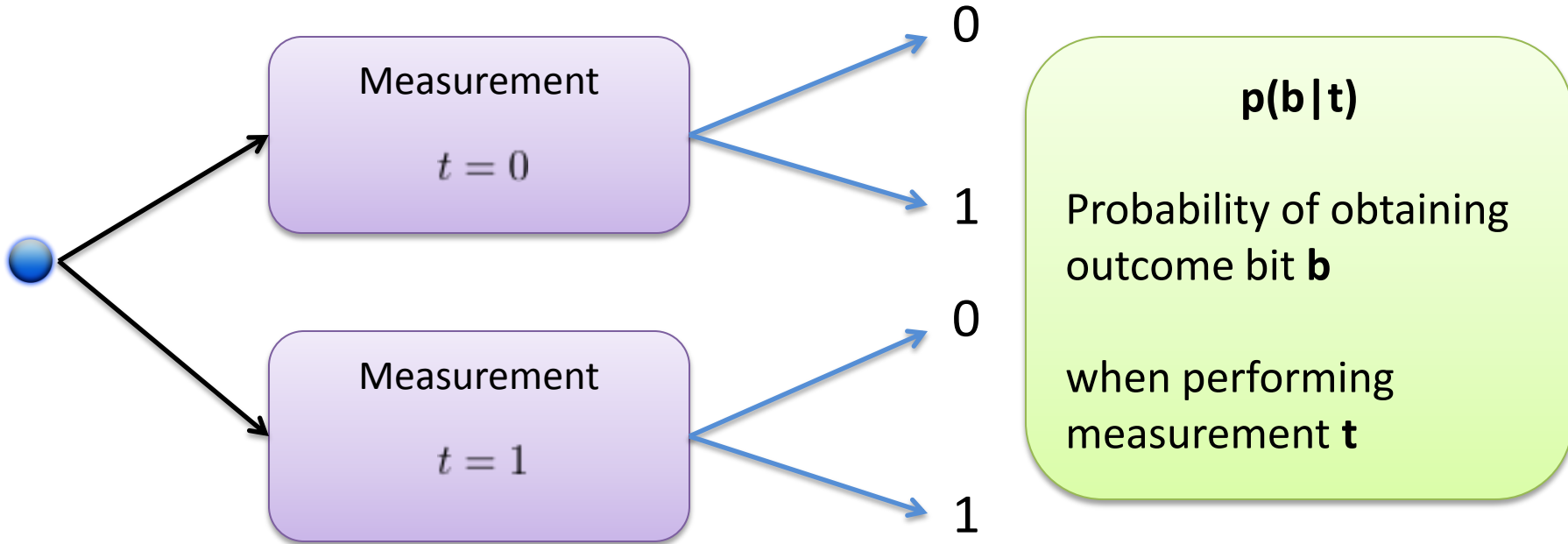
$$(0 < p < 1)$$

“Cannot know both the particle’s position and its momentum at the same time” observed by Werner Heisenberg in 1927.





# Uncertainty relations

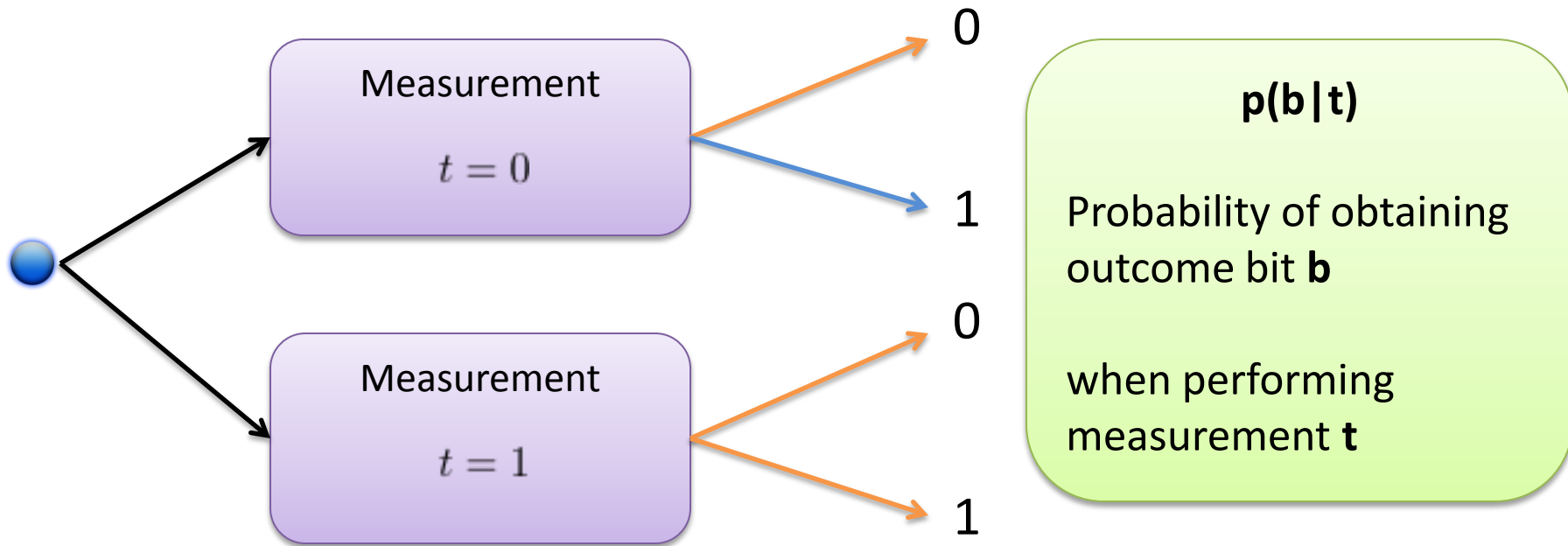


“Cannot know both the particle’s position and its momentum at the same time” observed by Werner Heisenberg in 1927.





# Uncertainty relations



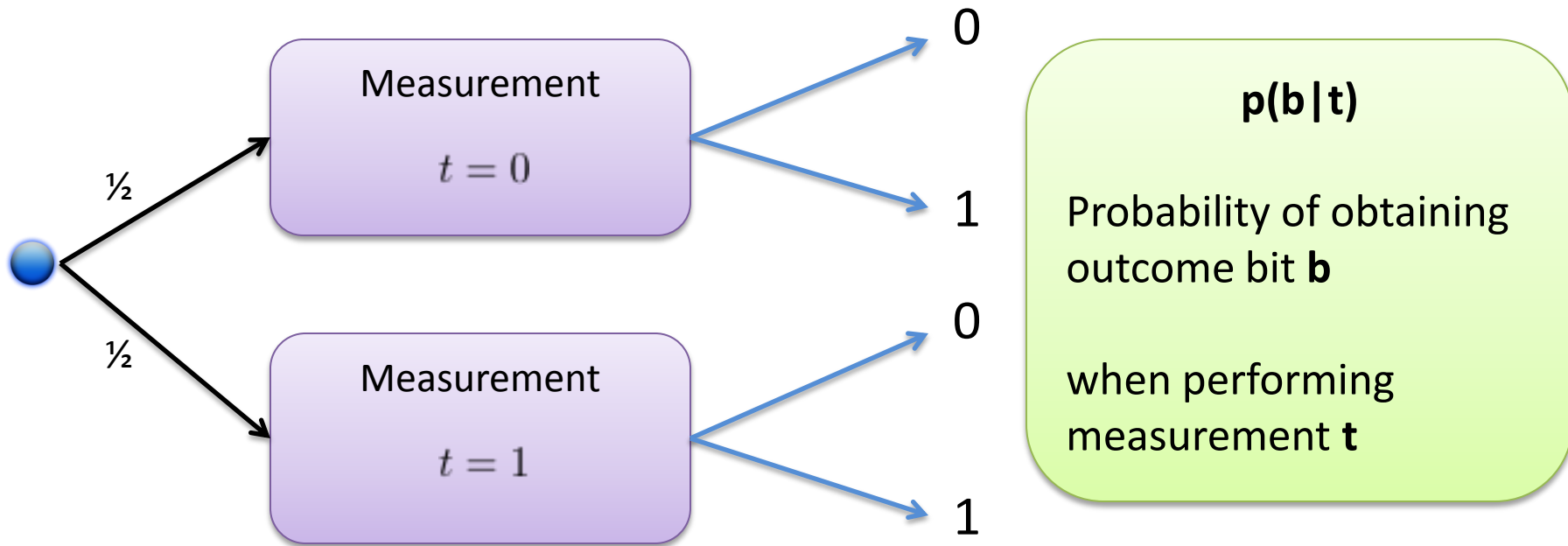
“Cannot know both the particle’s position and its momentum at the same time” observed by Werner Heisenberg in 1927.







# Choice of measurements

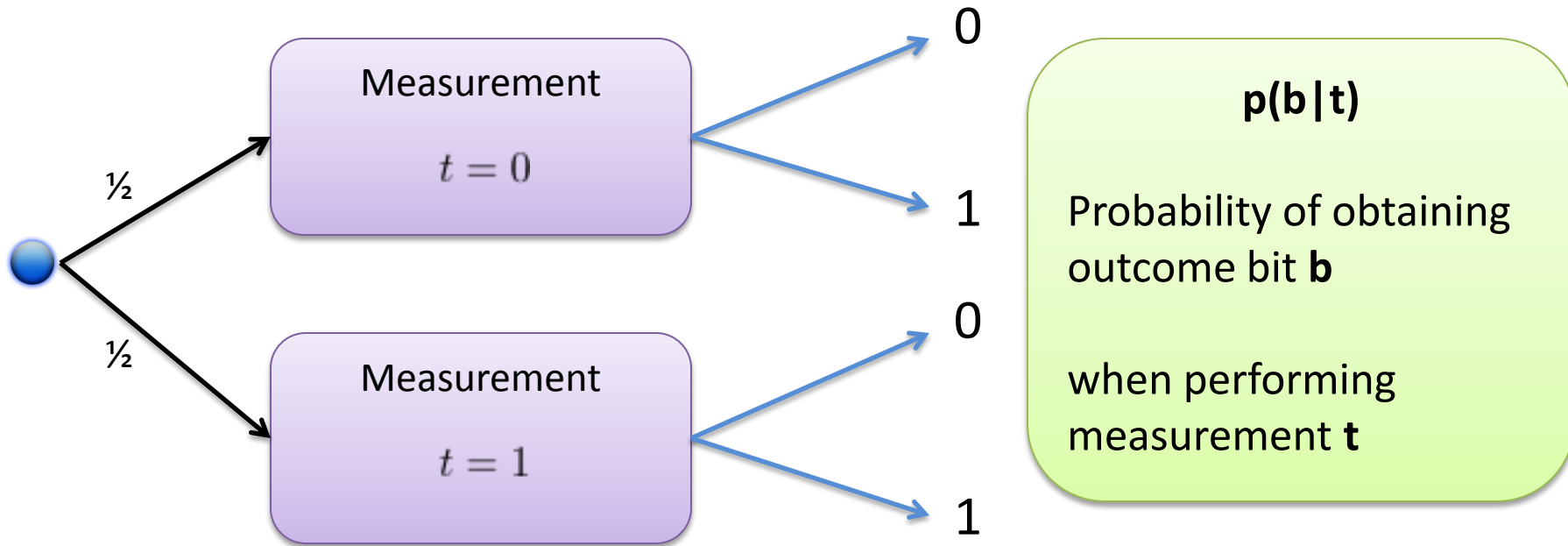


“Cannot know both the particle’s position and its momentum at the same time” observed by Werner Heisenberg in 1927.





# Fine-grained uncertainty relations



Suppose that for **any** quantum state:

$$\frac{1}{2} (p(0|t = 0) + p(0|t = 1)) \leq c$$

$$\frac{1}{2} (p(0|t = 0) + p(1|t = 1)) \leq c$$

$$\frac{1}{2} (p(1|t = 0) + p(0|t = 1)) \leq c$$

$$\frac{1}{2} (p(1|t = 0) + p(1|t = 1)) \leq c$$

**Fine-grained uncertainty relation:**

**Imagine  $c < 1$ .**

**Cannot have**

$$p(0|t = 0) = p(1|t = 1) = 1$$

**Cannot know both properties!**



# A change of perspective

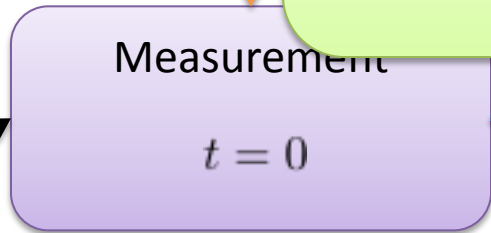
Probability of retrieving the first bit



$\frac{1}{2}$

$\frac{1}{2}$

Measurement



1

0

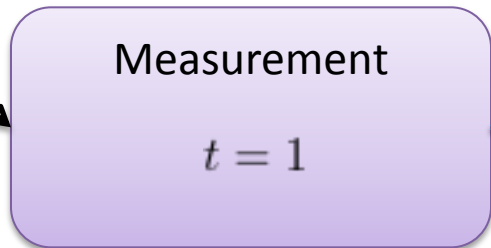
1

Uncertainty relation: "Cannot know both bits at the same time".

Encodes information in properties

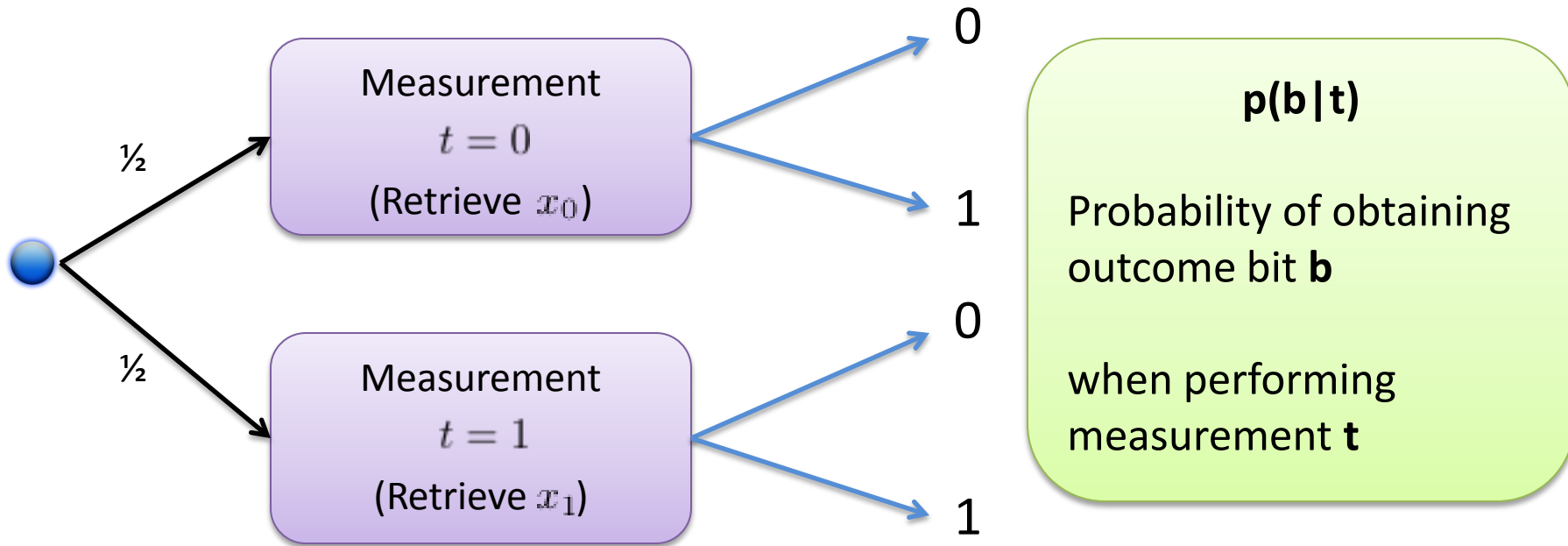
$$x_0, x_1 \in \{0, 1\}$$

Measurement to learn bit  $x_1$





# Encoding and decoding information



Suppose that for **any** quantum state:

$$\frac{1}{2} (p(0|t = 0) + p(0|t = 1)) \leq c$$

$$\frac{1}{2} (p(0|t = 0) + p(1|t = 1)) \leq c$$

$$\frac{1}{2} (p(1|t = 0) + p(0|t = 1)) \leq c$$

$$\frac{1}{2} (p(1|t = 0) + p(1|t = 1)) \leq c$$

## Coding interpretation:

Choose to retrieve  $x_0$  or  $x_1$  with probability  $\frac{1}{2}$

Average probability we correctly retrieve the desired bit from  $x_0 x_1 = (0, 0)$  for **any** encoding using those measurements.



# Maximally certain states

Suppose that for **any** quantum state:

$$\frac{1}{2} (p(0|t=0) + p(0|t=1)) \leq c$$

$$\frac{1}{2} (p(0|t=0) + p(1|t=1)) \leq c$$

$$\frac{1}{2} (p(1|t=0) + p(0|t=1)) \leq c$$

$$\frac{1}{2} (p(1|t=0) + p(1|t=1)) \leq c$$



# Maximally certain states

For the **maximally certain states**

$\frac{1}{2} (p(0 t=0) + p(0 t=1)) = c$	$\leftarrow \rho_{0,0}$	Encoding of $x = x_0x_1 = (0, 0)$
$\frac{1}{2} (p(0 t=0) + p(1 t=1)) = c$	$\leftarrow \rho_{0,1}$	Encoding of $x = x_0x_1 = (0, 1)$
$\frac{1}{2} (p(1 t=0) + p(0 t=1)) = c$	$\leftarrow \rho_{1,0}$	Encoding of $x = x_0x_1 = (1, 0)$
$\frac{1}{2} (p(1 t=0) + p(1 t=1)) = c$	$\leftarrow \rho_{1,1}$	Encoding of $x = x_0x_1 = (1, 1)$

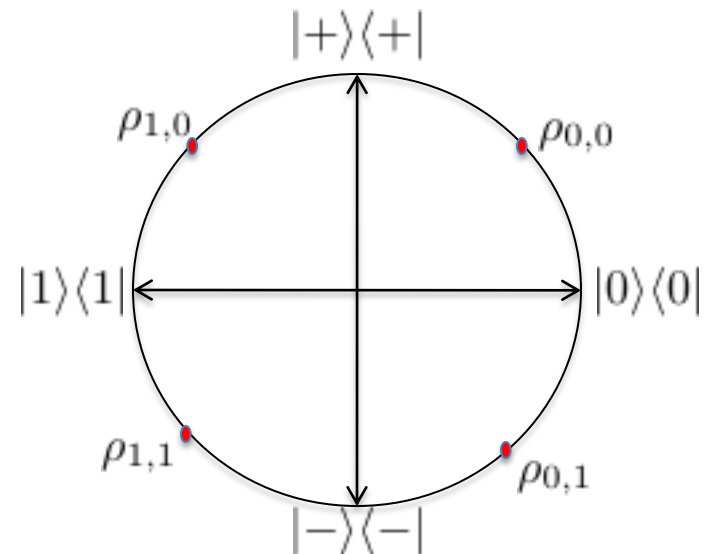
An example:

Measurements in Z and X basis

$$\frac{1}{2} \text{tr} [\rho (|0\rangle\langle 0| + |+\rangle\langle +|)] \leq \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

$$\frac{1}{2} \text{tr} [\rho (|0\rangle\langle 0| + |-\rangle\langle -|)] \leq \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

$\vdots$





# General form

Set of measurements  $T$

Set of outcomes  $B$

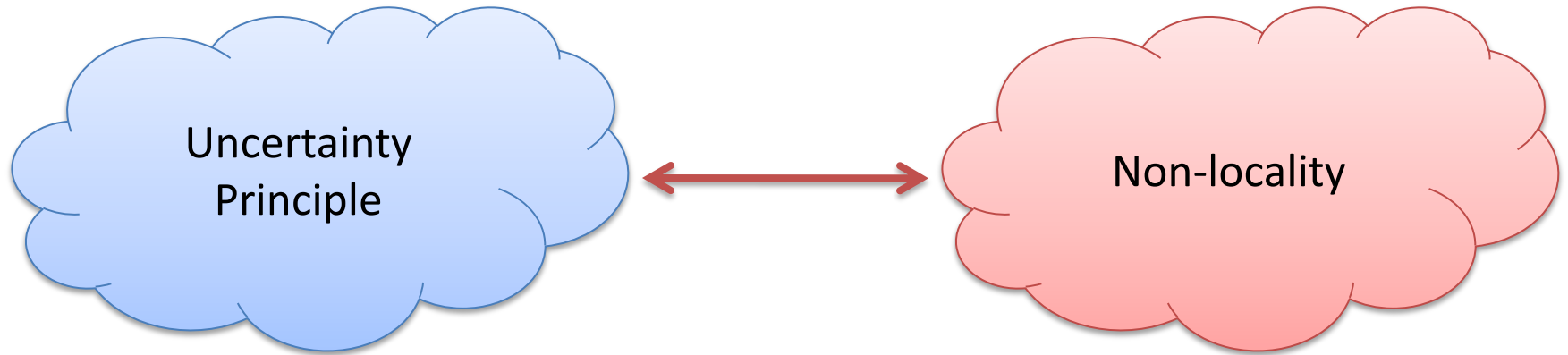
Probability distribution over measurements  $p(t)$

$$\mathcal{F}_{\text{UR}} = \left\{ \sum_t p(t)p(x_t|t) \leq \zeta_{\vec{x}} \mid \vec{x} = (x_1, x_2, \dots, x_{|T|}) \in B^{\times|T|} \right\}$$

Uncertainty relations have an operational interpretation telling us how well we can retrieve information using two particular measurements.

Maximally certain states are the best encoding with respect to those measurements.

# Linking two fundamental concepts

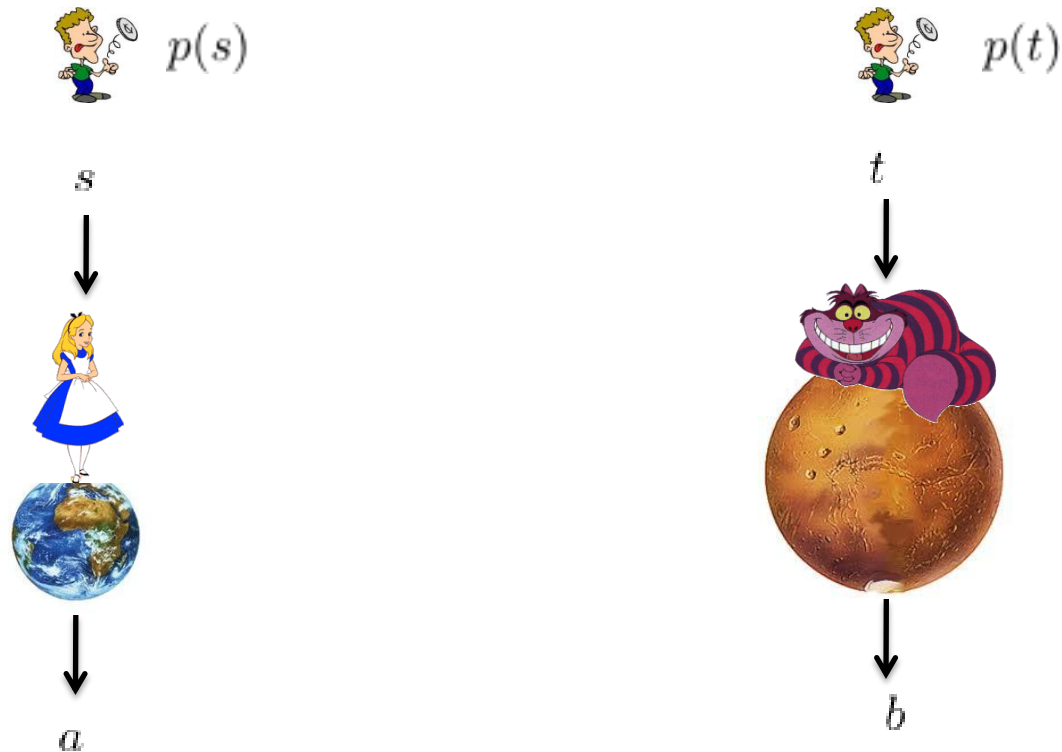


Steering





# Non-locality as a Game



Assumption about rules:

For all questions  $s, t$  and all answers  $a$  there exists exactly one winning answer  $b$  for Bob

Rules written as strings (W,Christandl, Doherty, '08):  $\vec{x}_{s,a} = (x_{s,a}^{t=0}, x_{s,a}^{t=1}, \dots)$

Correct answer for  $t = 0$       Correct answer for  $t = 1$



# Making Alice and Bob's life difficult..



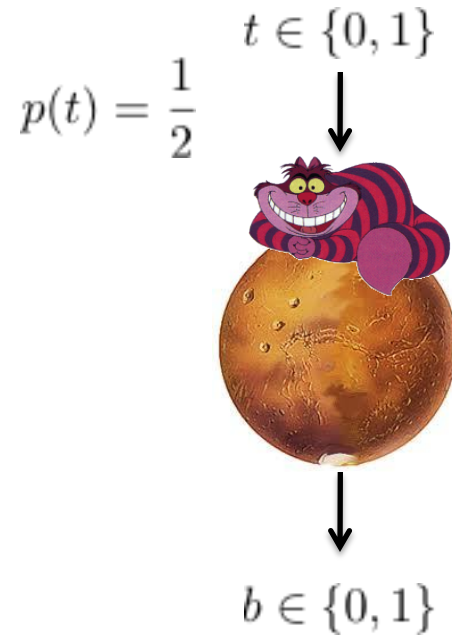
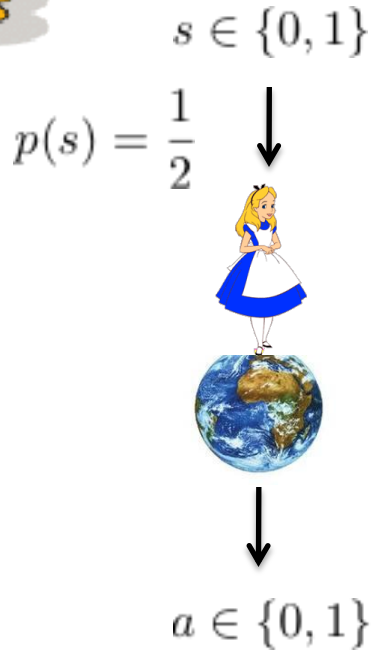
1. Can agree on any strategy beforehand
2. Once the game starts, they can no longer communicate

Amount of non-locality measured by winning probability

$$p_{\text{win}} = \sum_{s,t} p(s,t) \sum_{a,b} p(a,b = x_{s,a}^t | s,t)$$



# Example: CHSH as a game



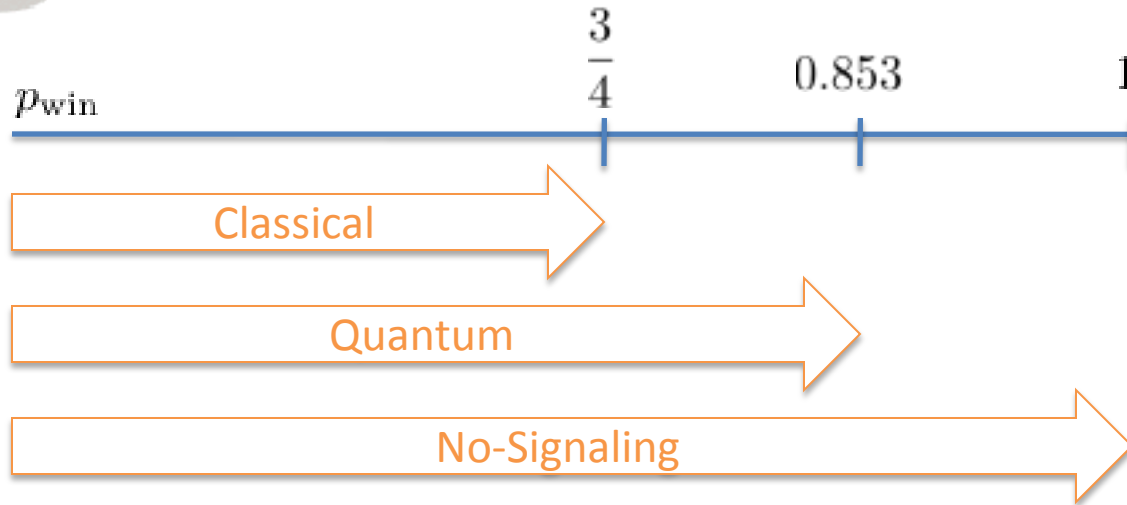
Rules as an equation  $s \cdot t = a + b \pmod{2}$

Rules as strings

$x_{0,0} = (0, 0)$	}	$s = 0$
$x_{0,1} = (1, 1)$		
$x_{1,0} = (0, 1)$	}	$s = 1$
$x_{1,1} = (1, 0)$		



# Why is non-locality limited?



Popescu and Rohrlich '96: Why is nature not more non-local?

More non-locality would allow for much better information processing:

Less communication (van Dam'00), store more information (verSteege, W '08),

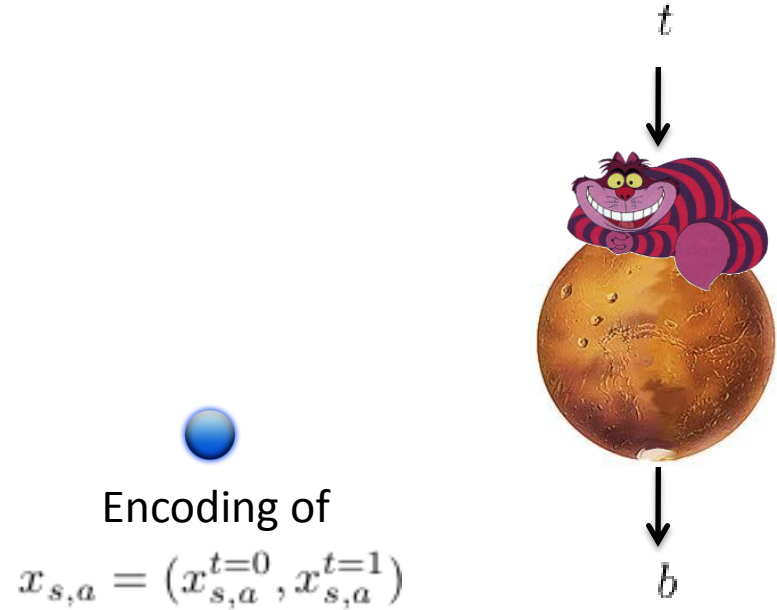
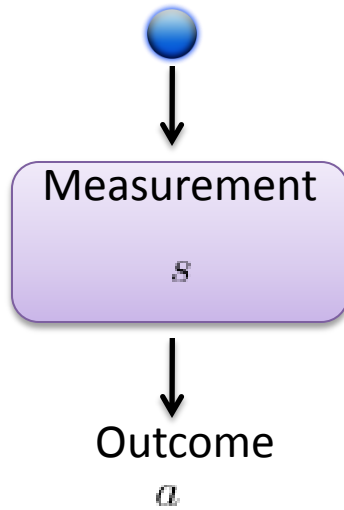
send more information per bit (Pawlowski, Paterek, Kaszlikowski, Scarani, Winter, Zukowski '09),...

Being locally quantum implies quantum correlations (Barnum, Beigi, Boixo, Elliot, Wehner '09)...

Would need to break the uncertainty principle for a set of measurements to have more non-locality.



# Producing answers



$$x_{0,0} = (0, 0)$$

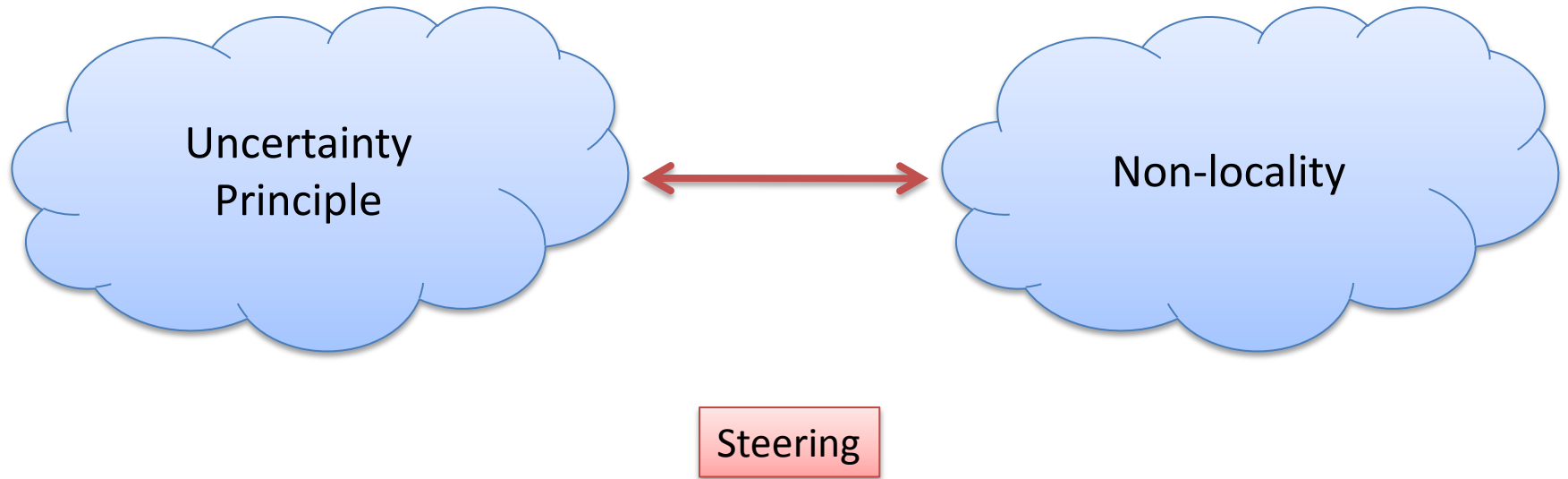
$$x_{0,1} = (1, 1)$$

$$x_{1,0} = (0, 1)$$

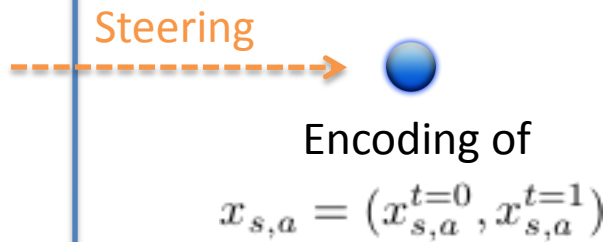
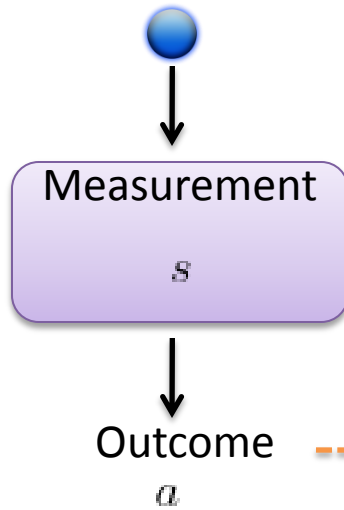
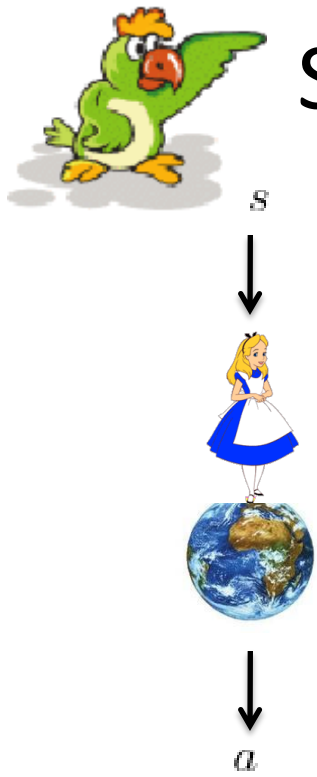
$$x_{1,1} = (1, 0)$$

To give the right answer Bob needs to retrieve bit  $x_{s,a}^t$  from the encoding

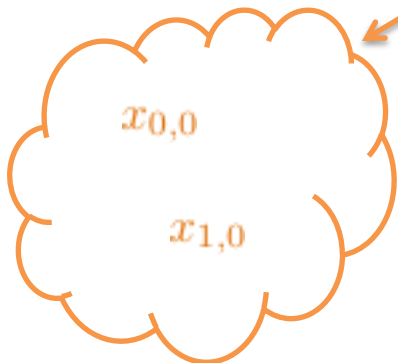
# Linking two fundamental concepts



# Steering (Schroedinger, 1935)



$$\mathcal{E}_{s=0} = \{p(a), \rho_{\vec{x}_{s=0,a}}\}_a$$



$s = 0$

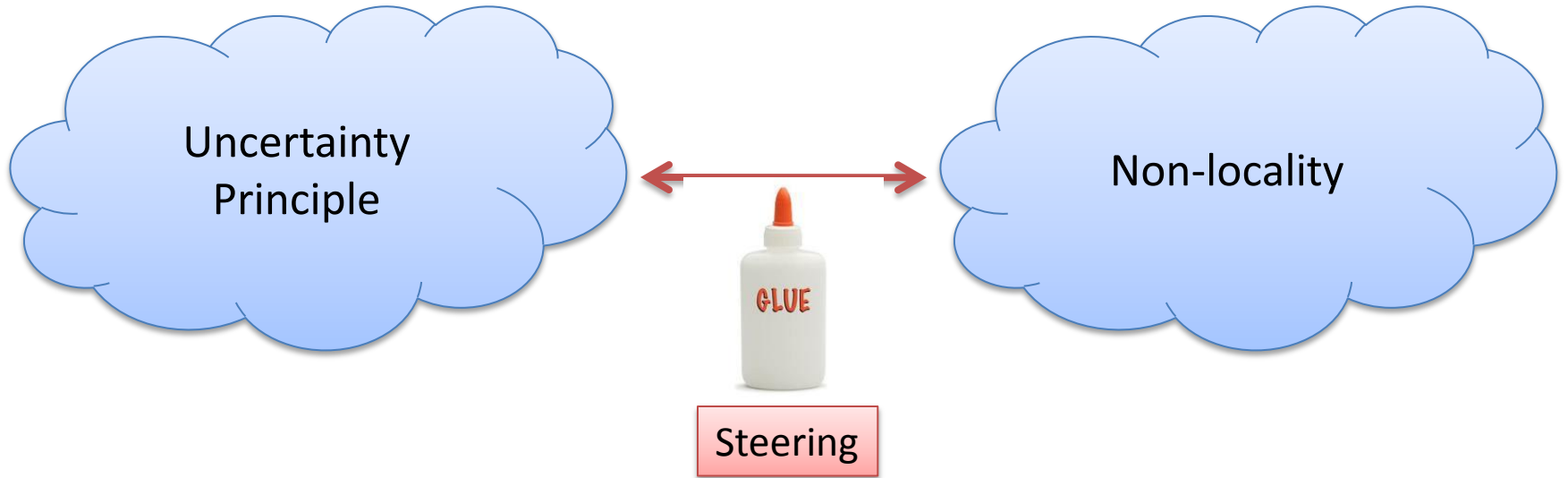


$s = 1$



$$\mathcal{E}_{s=1} = \{p(a), \rho_{\vec{x}_{s=1,a}}\}_a$$

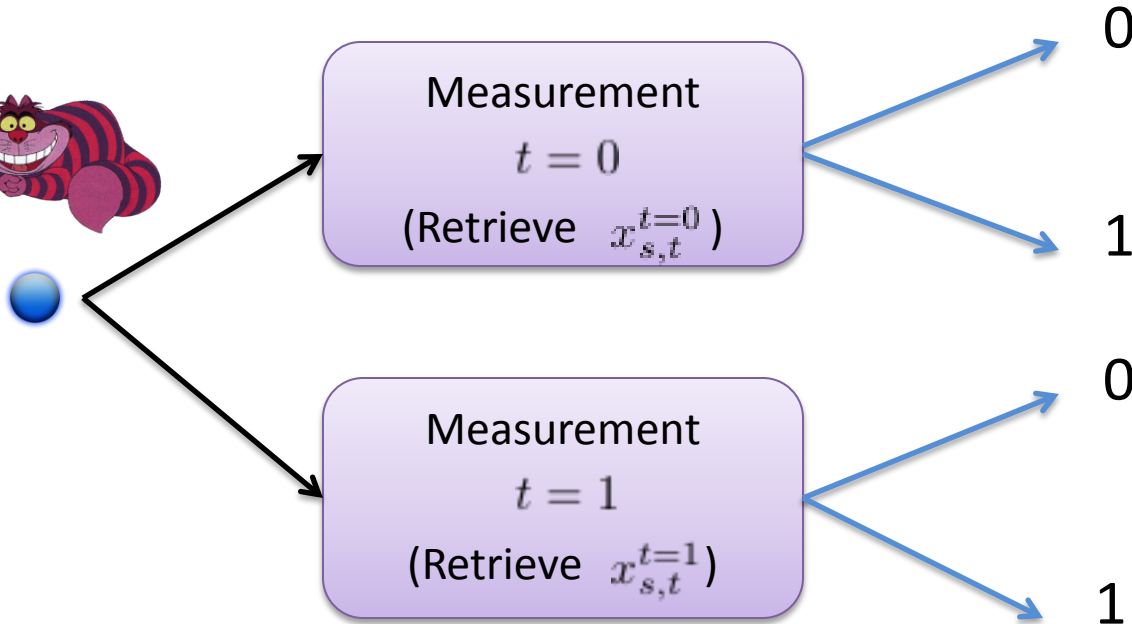
# Linking two fundamental concepts







# Producing answers



Uncertainty principle  
limits non-locality

$$\begin{aligned}
 p_{\text{win}} &= \sum_{s,t} p(s,t) \sum_{a,b} p(a, b = x_{s,a}^t | s, t) \\
 &= \sum_{s,a} p(s)p(a|s) \left[ \sum_t p(t)p(x_{s,a}^t | s, t, a) \right] \\
 &\leq \sum_{s,a} p(s)p(a|s) \zeta_{\vec{x}_{s,a}}
 \end{aligned}$$

Uncertainty relation, for any state

$$\sum_t p(t)p(x_{s,a}^t | t) \rho_{\vec{x}_{s,a}} \leq \zeta_{\vec{x}_{s,a}}$$



# Steering to the maximally certain states

1. Fix Bob's measurements
2. For these measurements, for the maximally certain states  $\rho_{\vec{x}_{s,a}}$

$$\sum_t p(t) p(x_{s,a}^t | t) = \zeta_{\vec{x}_{s,a}}$$

3. If Alice can steer to the maximally certain states

$$\begin{aligned} p_{\text{win}} &= \sum_{s,t} p(s,t) \sum_{a,b} p(a,b = x_{s,a}^t | s,t) \\ &= \sum_{s,a} p(s) p(a|s) \zeta_{\vec{x}_{s,a}} \end{aligned}$$

For XOR games, quantum Alice can steer to the maximally certain states of Bob's optimal measurements:

More non-locality would require a violation of the uncertainty principle.

The amount of non-locality is exactly determined by the amount of "uncertainty" and our ability to steer:

In any theory, going beyond the optimum winning probability would require a violation of the uncertainty principle with respect to the set of steerable states.



# Example CHSH

Rules

$$x_{0,0} = (0, 0)$$

$$x_{0,1} = (1, 1)$$

$$x_{1,0} = (0, 1)$$

$$x_{1,1} = (1, 0)$$

distribution over questions  $p(s) = p(t) = \frac{1}{2}$

Optimal measurements for Bob: measure Z and X.

$$\frac{1}{2} \text{tr} [\rho (|0\rangle\langle 0| + |+\rangle\langle +|)] \leq \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

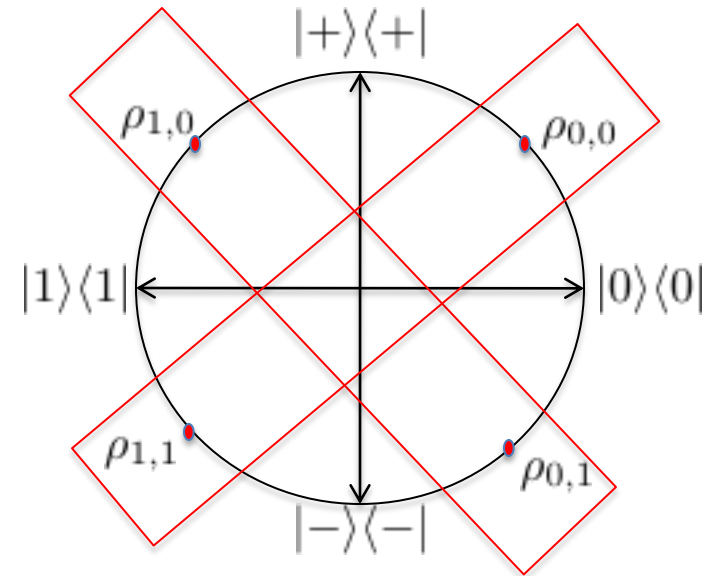
$$\frac{1}{2} \text{tr} [\rho (|0\rangle\langle 0| + |-\rangle\langle -|)] \leq \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

⋮

For the maximally certain states

$$\frac{1}{2} (\rho_{0,0} + \rho_{1,1}) = \frac{\text{id}}{2}$$

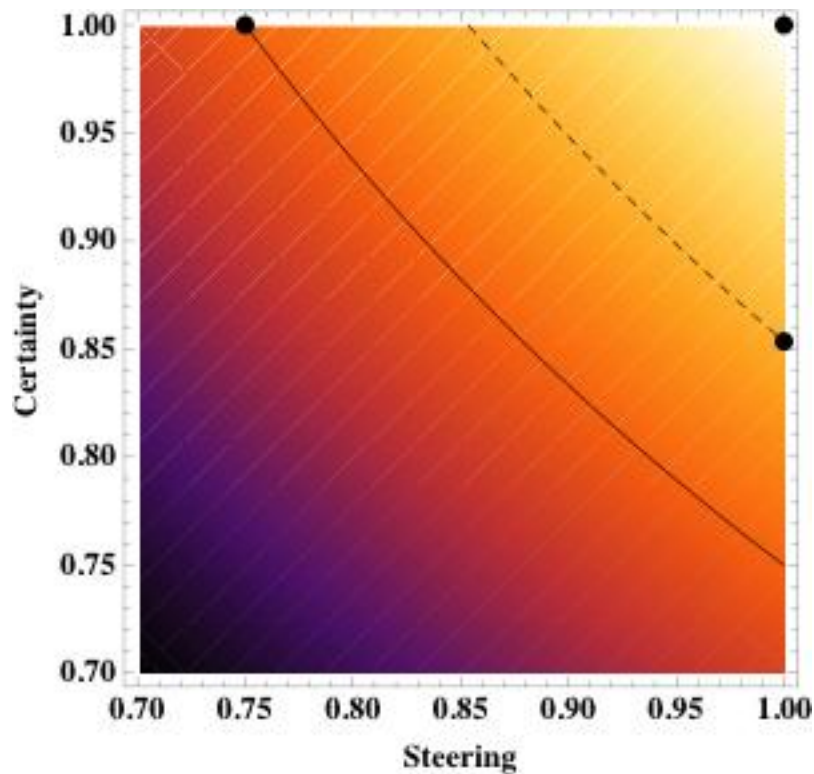
$$\frac{1}{2} (\rho_{1,0} + \rho_{0,1}) = \frac{\text{id}}{2}$$





# Non-locality for different theories

- Classical deterministic: No uncertainty, but also no steering.
- Quantum: Uncertainty, but perfect steering.
- No-signaling: No uncertainty, and perfect steering.

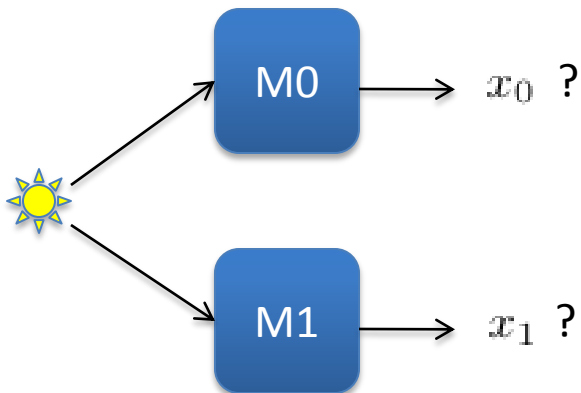


Uncertainty and non-locality are linked in **any** possible physical theory.



# Uncertainty vs. complementarity

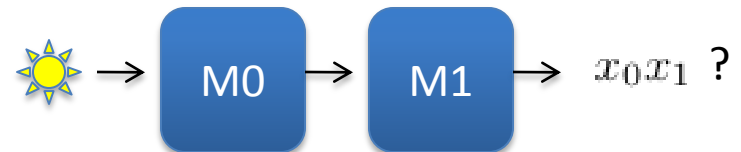
## Uncertainty



CHSH example:

$$\frac{1}{2} (p(x_0|t=0) + p(x_1|t=1)) \leq \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

## Complementarity



CHSH example: Cannot learn

$$x_0 \oplus x_1 = s$$

Could be equally uncertain (and non-local) but less complementary!

# Summary and open questions



- Linked two fundamental concepts
- Degree of non-locality is already determined by two concepts: uncertainty relations and steering
- Holds for *any* physical theory.





# Summary and open questions

- What properties fully characterize quantum theory?
- How strong can uncertainty relations for multiple measurements be in quantum theory?
- Uncertainty vs. complementarity?

Thank you!



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Postdoc Positions!