Zero-Knowledge Against Quantum Attacks

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January 16, 2006

Zero-Knowledge Proof Systems [Goldwasser, Micali & Rackoff, 1985]

Assume that a **promise problem** $A=(A_{\text{yes}},A_{\text{no}})$ has been fixed. A **zero-knowledge proof system** for the problem A is a pair (V,P) of interacting parties; a (computationally bounded) **verifier** and a **prover**.

Interaction:

Both parties receive an input string $x \in A_{\text{yes}} \cup A_{\text{no}}$, exchange messages with one another, and finally the verifier V produces an output string denoted (V,P)(x).

Conditions:

Completeness: If $x \in A_{yes}$, then it must be the case that (V,P)(x)=1 (accept) with high probability.

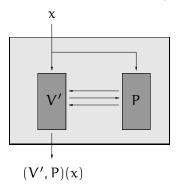
Soundness: If $x \in A_{no}$, then it must be the case that (V, P')(x) = 0 (reject) with high probability for every possible cheating prover P'.

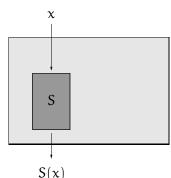
Zero-knowledge: If $x \in A_{yes}$, then no cheating verifier V' can extract knowledge from an interaction with P.

What does it mean to "extract knowledge"?

The notion of **knowledge** is a complexity-theoretic notion, and is different from **information**; it is formalized by means of the **simulator paradigm**.

Informally: a verifier V' learns nothing (i.e., fails to extract knowledge) from P if there exists a polynomial-time simulator S that produces an output that is **indistinguishable** from the output V' would produce when interacting with P on any $x \in A_{\text{yes}}$:

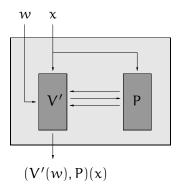


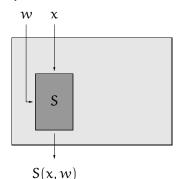


Auxiliary inputs

The previous informal definition is not quite strict enough to capture the notion of zero-knowledge, and gives rise to a class of protocols lacking certain desirable properties...

We need to allow the cheating verifier V' (as well as the simulator S) to take an **auxiliary input** string w. The outputs of these two processes should be indistinguishable provided $x \in A_{yes}$:





Auxiliary inputs

This **auxiliary input** definition captures the idea that zero-knowledge proofs should not **increase** knowledge, and is closed under sequential composition.

Definition of Zero-Knowledge (classical)

An interactive proof system (P,V) for a given problem $A=(A_{yes},A_{no})$ is **zero-knowledge** if, for every polynomial-time verifier V' there exists a **polynomial-time simulator** S such that, for every w and $x \in A_{yes}$,

$$(V'(w), P)(x)$$
 and $S(x, w)$

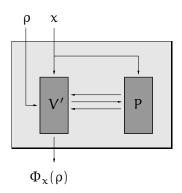
are indistinguishable.* [Goldwasser, Micali & Rackoff, 1989].

^{*} Different notions of indistinguishability give rise to different variants of zero-knowledge, such as **statistical** and **computational** zero-knowledge.

Quantum version of the definition

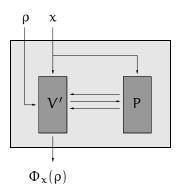
Suppose that some verifier V' tries to use **quantum information** to extract knowledge from P. (Note that the prover P is still classical, so the input x and any information exchanged between V' and P must be classical.)

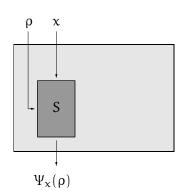
The interaction between V' and P on input x induces some **admissible mapping** on the auxiliary input:



Quantum version of the definition

If P is zero-knowledge even against a verifier V' that uses quantum information, then there should exist a simulator S that performs an admissible mapping Ψ_x on the auxiliary input that is **indistinguishable** from Φ_x (when $x \in A_{\text{yes}}$):





Problem with the quantum definition?

These definitions are fairly straightforward...but have been considered problematic for several years. (The problem was apparently first identified by Jeroen van de Graaf in his 1997 PhD thesis.)

The problem: No nontrivial protocols were previously shown to be zero-knowledge with respect to these definitions, even protocols already proved zero-knowledge in the classical setting.

In order to describe the problem, it will be helpful to consider a simple and well-known zero-knowledge proof system for the **Graph Isomorphism** problem:

Input: Two graphs G_0 and G_1 (given by adjacency matrices).

Yes: G_0 and G_1 are isomorphic ($G_0 \cong G_1$).

No: G_0 and G_1 are not isomorphic ($G_0 \not\cong G_1$).

A zero-knowledge proof system for Graph Isomorphism

The following protocol (described for honest parties) is a zero-knowledge protocol for Graph Isomorphism [Goldreich, Micali & Widgerson, 1991].

The GMW Graph Isomorphism Protocol

Assume the input is a pair (G_0,G_1) of \mathfrak{n} -vertex graphs. Let $\sigma\in S_\mathfrak{n}$ be a permutation satisfying $\sigma(G_1)=G_0$ if $G_0\cong G_1$, and let σ be arbitrary otherwise.

Prover's step 1: Choose $\pi \in S_n$ uniformly at random and send $H = \pi(G_0)$ to the verifier.

Verifier's step 1: Choose $\alpha \in \{0,1\}$ randomly and send α to the prover. (Implicit: challenge prover to show $H \cong G_{\alpha}$.)

Prover's step 2: Let $\tau = \pi \sigma^{\alpha}$ and send τ to the verifier.

Verifier's step 2: Accept if $\tau(G_{\alpha}) = H$, reject otherwise.

Sequential repetition reduces soundness error...

Zero-knowledge property for the GMW protocol

The **completeness** and **soundness** properties are straightforward. Let us consider the **zero-knowledge** property...

Consider a **classical** cheating verifier V':

Verifier's step 1: Perform some **arbitrary** polynomial-time computation on (G_0, G_1) , auxiliary input w, and H to obtain $a \in \{0, 1\}$. Send a to P.

Verifier's step 2: Perform some **arbitrary** polynomial-time computation on (G_0, G_1) , auxiliary input w, H, and τ to produce output.

Simulator for V':

- 1. Choose $b \in \{0,1\}$ and $\tau \in S_{\pi}$ uniformly, and let $H = \tau(G_b)$.
- 2. Simulate whatever V^\prime does given prover message H. Let α denote the resulting message back to the prover.
- 3. If $a \neq b$ then rewind: go back to step 1 and try again.
- 4. Output whatever V' would after receiving τ .

Simulator for a cheating quantum verifier?

Suppose that we have a cheating **quantum** verifier V' that starts the protocol with an auxiliary quantum register W.

Verifier's step 1: Perform some arbitrary polynomial-time quantum computation on (G_0, G_1) , auxiliary input register W, and H to obtain $a \in \{0, 1\}$. Send a to P.

For example: let α be the outcome of some binary-valued projective measurement $\{\Pi_0^H, \Pi_1^H\}$ of W that depends on H.

Verifier's step 2: Perform some arbitrary polynomial-time quantum computation to produce an output.

How can we simulate such a verifier?

The "no quantum rewinding" issue

Two principles are working against us:

- The no cloning theorem prevents making a copy of the auxiliary input register's state.
- Measurements are irreversible.

Suppose that we randomly choose b and τ , and let $H=\tau(G_b)$ as for our simulator before. If the simulator guesses incorrectly (meaning $a\neq b$), then the original state of W may not be recoverable.

"Rewinding by reversing the unitary transformation induced by [the verifier], or taking snapshots is impossible.

But... showing that rewinding by reversing or by taking snapshots is impossible does not show that no other ways to rewind in polynomial time exist."

[VAN DE GRAAF, 1997]

New results

In the remainder of this talk I will argue that the GMW Graph Isomorphism protocol is indeed zero-knowledge against quantum verifiers:

- For any quantum verifier V', there exists a simulator S that induces
 precisely the same admissible mapping as the interaction between V'
 and P (on a "yes" input to the problem).
- The method gives a way to "rewind" the simulator, but it requires more than just reversing the verifier's actions. (The entire simulation will be quantum, even though the prover is classical.)
- The method generalizes to several other protocols (but I will only discuss the Graph Isomorphism example in this talk for simplicity).

Assumptions on V'

Assume V' uses three registers:

W: stores the auxiliary input.

V: represents workspace of arbitrary size.

A: single qubit representing the message sent by V'.

Register W starts in the auxiliary state, and registers V and A are initialized to all zeroes.

Assume V' operates as follows:

- For each graph H on n vertices, V' has a corresponding unitary transformation V_H that acts on (W, V, A).
- Upon receiving H from P, the V' applies V_H to (W, V, A), measures
 A in the standard basis, and sends the result a to P.
- After P responds with some permutation τ , V' simply outputs (W,V,A) along with the prover messages H and τ .

Simulator construction

The simulator will use registers W, V, and A along with:

 P_1 : stores the prover's first message.

B: stores the simulator's guess b for a.

P₂: stores the prover's second message.

R: stores "randomness" used to generate transcripts.

Define a unitary operator V on (W, V, A, P_1) that represents a unitary realization of V':

$$V = \sum_{H} V_{H} \otimes \left| H \right\rangle \left\langle H \right|.$$

Define T to be a unitary operation on registers (P₁, B, P₂, R) for which

$$T: |00\cdots 0\rangle \mapsto \frac{1}{\sqrt{2n!}} \sum_{b,\tau} |\tau(G_b)\rangle |b\rangle |\tau\rangle |b,\tau\rangle.$$

The operation T produces a superposition over *transcripts*.

Simulator construction

Now define the simulator as follows:

Simulator

- 1. Perform T, followed by V.
- 2. Perform a measurement $\{\Pi_0, \Pi_1\}$ whose outcome corresponds to the XOR of A and B (in the computational basis).
- 3. If the measurement outcome is 1, we need to rewind and try again:
 - Perform V* followed by T*.
 - Perform a **phase flip** in case any of the qubits in any of the registers (V, A, P₁, B, P₂, R) is set to 1 (i.e., perform $2\Delta I$, where $\Delta = I_W \otimes |00 \cdots 0\rangle \langle 00 \cdots 0|$.)
 - Perform T followed by V.
- 4. Output registers (W, V, A, P₁, P₂). (Registers B and R are traced out.)

Analysis of simulator

Assume that the auxiliary input is $|\psi\rangle$, and $x=(G_0,G_1)$ for $G_0\cong G_1.$ Let

$$|\phi\rangle = |\psi\rangle |00 \cdots 0\rangle$$

be the state of all registers given this input.

The simulator performs T, then V, then measures w.r.t. $\{\Pi_0, \Pi_1\}$. Assuming $G_0 \cong G_1$, the outcome will always be uniformly distributed.

First, suppose that the measurement $\{\Pi_0, \Pi_1\}$ gives **outcome 0**. The resulting state of all registers is

$$|\sigma_0\rangle = \sqrt{2}\Pi_0 V T |\phi\rangle \,. \label{eq:sigma0}$$

This is the target state: it represents a successful simulation because

$$\operatorname{tr}_{\mathsf{B},\mathsf{R}}\left|\sigma_{0}\right\rangle\left\langle\sigma_{0}\right|=\Phi(\left|\psi\right\rangle\left\langle\psi\right|).$$

(Nothing is surprising here... the simulator has been lucky and didn't need to rewind.)

Analysis of simulator

Suppose on the other hand that the **measurement outcome was 1**. The resulting state is

$$|\sigma_1\rangle = \sqrt{2}\Pi_1 V T |\phi\rangle$$
 .

Time to rewind and try again...

Performing the "rewind and try again" procedure results in the state

$$VT(2\Delta-I)T^*V^*\left|\sigma_1\right\rangle.$$

Claim

$$VT(2\Delta-I)T^*V^*|\sigma_1\rangle=|\sigma_0\rangle$$
 (the target state).

Note: this would not happen for **arbitrary** choices of $|\phi\rangle$, V, T, Π_0 , Π_1 , etc. . . the claim relies on the fact that the measurement $\{\Pi_0,\Pi_1\}$ gives outcome 0 and 1 with equal probability for **all** choices of $|\psi\rangle$.

Proof of claim

The fact that the measurement $\{\Pi_0,\Pi_1\}$ gives outcomes 0 and 1 with equal probability for **all** choice of $|\psi\rangle$ implies

$$\Delta T^*V^*\Pi_0VT\Delta = \Delta T^*V^*\Pi_1VT\Delta = \frac{1}{2}\Delta.$$

Therefore

$$\begin{split} \langle \sigma_0 | VT(2\Delta-I)T^*V^* | \sigma_1 \rangle \\ &= 2 \left\langle \phi | T^*V^*\Pi_0 VT(2\Delta-I)T^*V^*\Pi_1 VT | \phi \right\rangle \\ &= 4 \left\langle \phi | T^*V^*\Pi_0 VT\Delta T^*V^*\Pi_1 VT | \phi \right\rangle \\ &- 2 \left\langle \phi | T^*V^*\Pi_0 VTT^*V^*\Pi_1 VT | \phi \right\rangle \\ &= 4 \left\langle \phi | \Delta T^*V^*\Pi_0 VT\Delta T^*V^*\Pi_1 VT\Delta | \phi \right\rangle \\ &= \left\langle \phi | \Delta | \phi \right\rangle \\ &= 1, \end{split}$$

so
$$VT(2\Delta - I)T^*V^*|\sigma_1\rangle = |\sigma_0\rangle$$
.

Analysis of simulator

This establishes that the admissible map Ψ agrees with the map Φ corresponding to the actual interaction on all pure state auxiliary inputs:

$$\Psi(\ket{\psi}\bra{\psi}) = \Phi(\ket{\psi}\bra{\psi})$$

for all $|\psi\rangle$.

Admissible maps are **completely determined** by their actions on pure state inputs, however, so

$$\Psi = \Phi$$
;

the simulator **agrees precisely** with the actual interaction on **every possible state** of the auxiliary input register (including the possibility it is entangled with another register).

Other protocols

The simulation method just described can be adapted to prove several other protocols are zero-knowledge against quantum attacks, including:

 Quantum protocols for any problem having an honest verifier quantum statistical zero-knowledge proof system:

$$QSZK = QSZK_{HV}$$
.

- The Goldreich-Micali-Wigderson Graph 3-Coloring protocol assuming unconditionally binding and quantum computationally concealing bit commitments. (See [ADCOCK & CLEVE, 2002].)
- Presumably several other proof systems...

Adapting the simulator to other protocols may require iterating the "rewind and try again" process.

Future work/open questions

- 1. Find further applications and generalizations of the method.
- 2. Identify limitations of the method.
- 3. Identify good candidates for quantum one-way permutations.